

Killing constrained supersymmetry softly...

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Thanks for the invitation to UC Irvine!



- ▶ Physics beyond the Standard Model at the LHC?
- ▶ Great expectations: the pre-LHC gold rush
- ▶ Global fits of supersymmetric models
- ▶ Beyond supersymmetry: simplified models
- ▶ Conclusions

Why new physics at the LHC?

- ▶ The naturalness problem: why is $M_{\text{Higgs}} \ll M_{\text{Planck}}$?



$$\delta M_{\text{H}}^2 \sim \frac{3\lambda_t^2}{8\pi^2} \Lambda_{\text{UV}}^2 \sim (0.3 \Lambda_{\text{UV}})^2$$

→ new colored (top) partners with mass below about 500 GeV?

- ▶ Dark matter



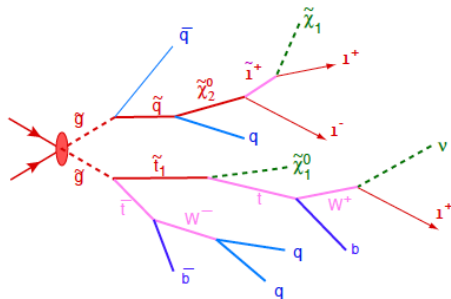
weakly interacting massive particle with mass below about 1 TeV?

New physics at the LHC?

► Naturalness & dark matter

→ new particles at the TeV-scale including a stable WIMP

→ generic LHC new physics signature: cascade decays with $E_{T,miss}$



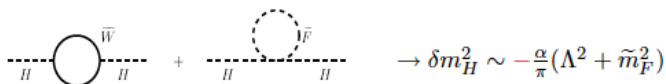
Supersymmetry is a prime example for this class of models.

Symmetry between fermions and bosons: $Q|\text{boson}\rangle = |\text{fermion}\rangle$
 $Q|\text{fermion}\rangle = |\text{boson}\rangle$
with algebra $\{Q_\alpha, Q_\beta^\dagger\} = (\sigma^\mu)_{\alpha\beta} P_\mu$

→ SUSY is the unique maximal external symmetry in nature.

→ One needs to introduce superpartners to the Standard Model particles.

✓ SUSY protects the Higgs mass from large radiative corrections:


$$\text{---} \text{H} \text{---} \text{---} \text{H} \text{---} + \text{---} \text{H} \text{---} \text{---} \text{H} \text{---} \rightarrow \delta m_H^2 \sim -\frac{\alpha}{\pi}(\Lambda^2 + \tilde{m}_F^2)$$

$$\delta m_H^2 \sim \frac{\alpha}{\pi}(m_F^2 - \tilde{m}_F^2) \rightarrow \text{no fine-tuning if } \tilde{m} \lesssim \mathcal{O}(1 \text{ TeV})$$

✓ SUSY allows for coupling unification, radiative EWSB,
dark matter (assuming R-parity), ...

The Minimal Supersymmetric extension of the SM

- ▶ **external symmetries:** Poincaré symmetry & supersymmetry
- ▶ **internal symmetries:** $SU(3) \otimes SU(2) \otimes U(1)$ gauge symmetries

Gauge Bosons $S = 1$ gluon, W^\pm, Z, γ	Gauginos $S = 1/2$ gluino, $\tilde{W}, \tilde{Z}, \tilde{\gamma}$
Fermions $S = 1/2$ $\begin{pmatrix} u_L \\ d_L \end{pmatrix} \begin{pmatrix} \nu_L^e \\ e_L \end{pmatrix}$ u_R, d_R, e_R	Sfermions $S = 0$ $\begin{pmatrix} \tilde{u}_L \\ \tilde{d}_L \end{pmatrix} \begin{pmatrix} \tilde{\nu}_L^e \\ \tilde{e}_L \end{pmatrix}$ $\tilde{u}_R, \tilde{d}_R, \tilde{e}_R$
Higgs $\begin{pmatrix} H_2^0 \\ H_2^- \end{pmatrix} \begin{pmatrix} H_1^+ \\ H_1^0 \end{pmatrix}$	Higgsinos $\begin{pmatrix} \tilde{H}_2^0 \\ \tilde{H}_2^- \end{pmatrix} \begin{pmatrix} \tilde{H}_1^+ \\ \tilde{H}_1^0 \end{pmatrix}$

- ▶ SUSY breaking \rightarrow sparticle masses?

Searches for supersymmetry

The SUSY parameter space is strongly constrained by

- ▶ SM precision observables:

$\text{BR}(b \rightarrow s\gamma)$, $\text{BR}(B_s \rightarrow \mu\mu)$, $\text{BR}(b \rightarrow \tau\nu)$, Δm_{B_s} , $(g-2)_\mu$, m_W , $\sin^2 \theta_{\text{eff}}$

- ▶ astrophysical observations:

Ω_{DM} , direct and indirect DM detection limits

- ▶ direct sparticle and Higgs boson search limits from colliders:

$m_{\tilde{\chi}^\pm}$, LEP limits on MSSM Higgs bosons

- ▶ LHC SUSY exclusions from jets+ $E_{T_{\text{miss}}}$ searches

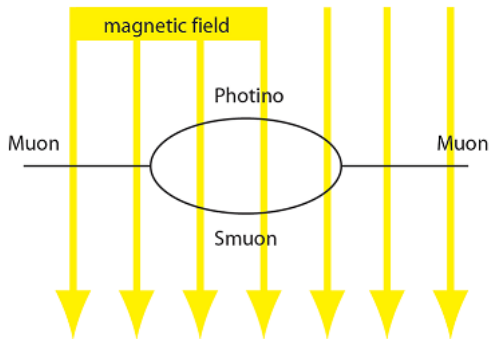
- ▶ the LHC Higgs signal

→ test supersymmetric models through global fits

[see e.g. SFitter, Mastercode, SuperBayes and Fittino]

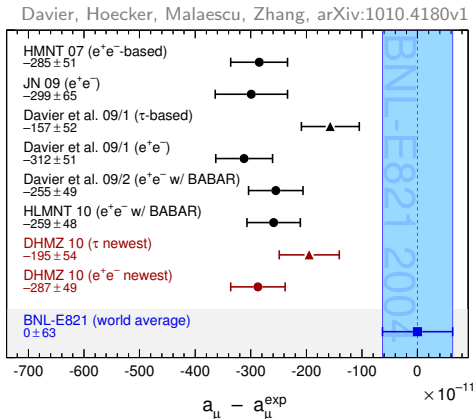
Indirect SUSY searches

- ▶ the anomalous magnetic moment of the muon $(g - 2)_\mu$:



→ SUSY loops: $a_\mu^{\text{SUSY}} \sim \text{sgn}(\mu) \tan\beta M_{\text{SUSY}}^{-2}$

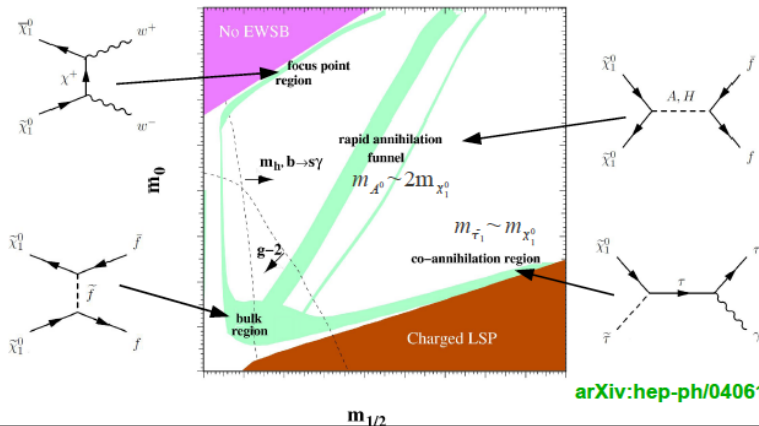
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Indirect SUSY searches

- Ω_{DM} is too large for large parts of the MSSM parameter space, special annihilation mechanisms are needed:



The constrained Minimal SuperSymmetric Model

Current data cannot constrain general SUSY models

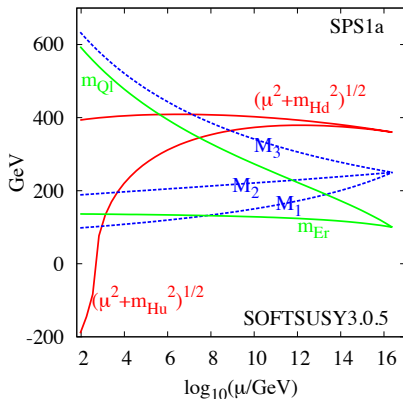
→ Consider the **constrained MSSM**: a model with universal scalar and fermion sparticle masses M_0 and $M_{1/2}$ at the GUT scale

The constrained Minimal SuperSymmetric Model

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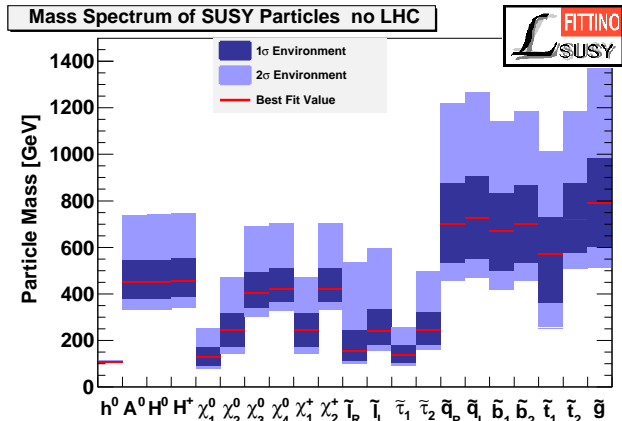
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low scale masses ← RG evolution → high scale parameters



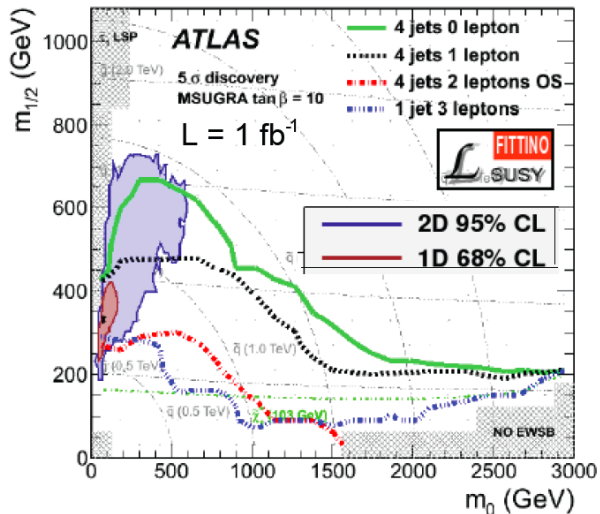
The constrained MSSM: global fits pre-LHC

cMSSM fit to B , K and EWK observables, $(g - 2)_\mu$ and Ω_{DM} :



→ SUSY is just around the corner...

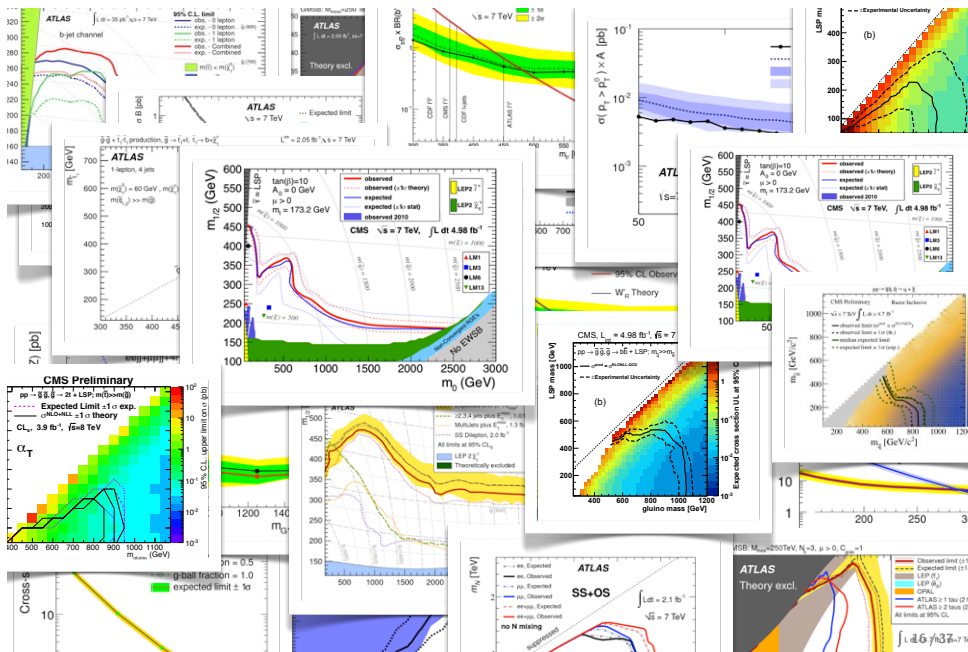
The constrained MSSM: great expectations



SUSY! SUSY! SUSY from the LHC!



Summary of BSM searches at the LHC: limits, limits and more limits...



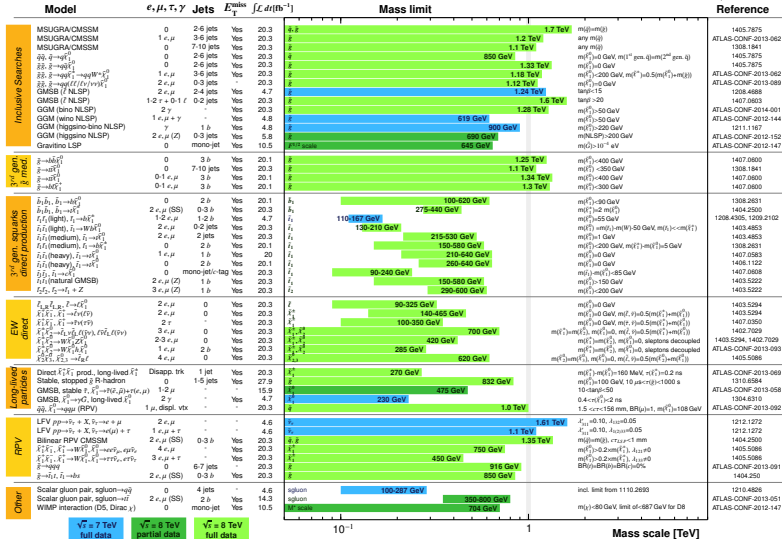
Summary of SUSY searches at the LHC

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: ICHEP 2014

ATLAS Preliminary

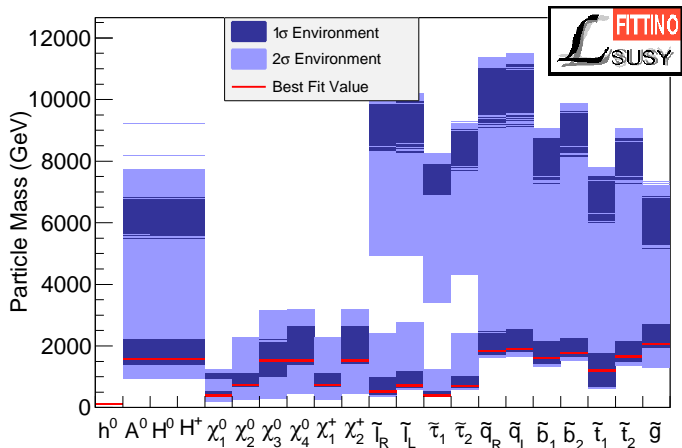
$\sqrt{s} = 7, 8 \text{ TeV}$



*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

The constrained MSSM: global fits post LHC-8

cMSSM fit including LHC exclusions and Higgs properties:



The constrained MSSM: global fits post LHC-8

To summarize the picture in the constrained MSSM:

- ▶ sparticles and H, A, H^\pm beyond the current LHC reach
 - ▶ branching ratios of the light Higgs h close to the SM values
 - ▶ the branching ratio for $B_s \rightarrow \mu\mu$ close to the SM value
 - ▶ no dark matter signal in current direct or indirect searches
-
- ▶ The cMSSM looks like the SM with dark matter.
 - ▶ The cMSSM could mean grim prospects for LHC phenomenology.
 - ▶ The cMSSM cannot solve the hierarchy problem.

We find that the cMSSM is pretty dull. But is it already dead?

How dead is the cMSSM?

Estimate the *p*-value:

If the cMSSM at the best fit point is the true description of nature, what is the probability p to get a minimal χ^2 as bad or worse than the one observed?

For a set of observables with Gaussian uncertainties, only need # of degrees of freedom and χ_{\min} to calculate the *p*-value.

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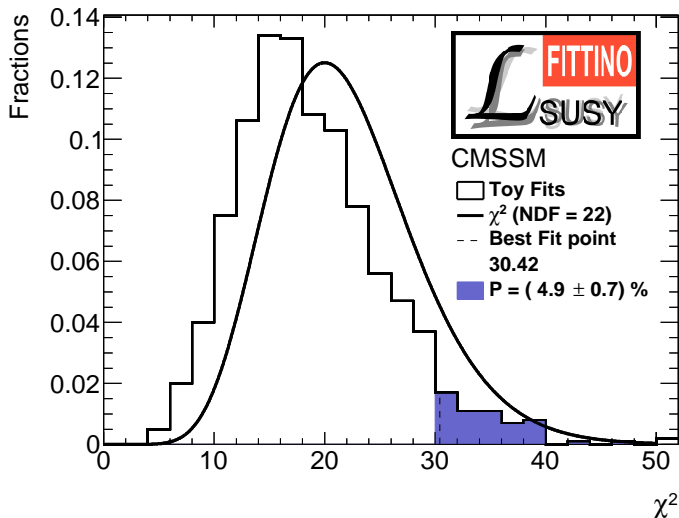
For a set of observables with Gaussian uncertainties, only need # of degrees of freedom and χ_{\min} to calculate the *p*-value.

Unfortunately, the uncertainties are, in general, non-Gaussian, so we have to perform *toy fits*:

- ▶ generate (1000) pseudo-measurements by smearing the experimental observables about the best-fit prediction;
- ▶ repeat the global fit for each of these pseudo-measurements.

How often do we get a minimal χ^2 as bad or worse than the one of the best-fit point? → *p*-value

How dead is the cMSSM?



How dead is the cMSSM? The impact of $(g - 2)_\mu$ and the Higgs

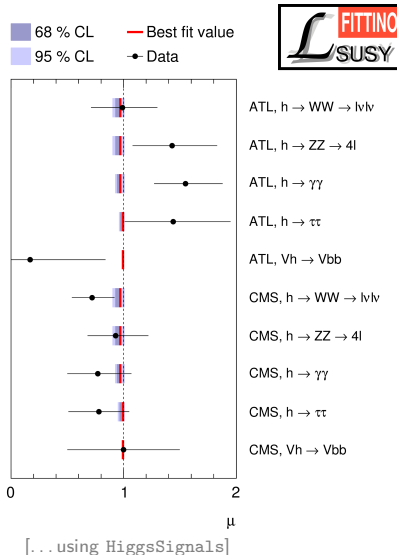
- ▶ without $(g - 2)_\mu$:
 $p = (51 \pm 3)\%$

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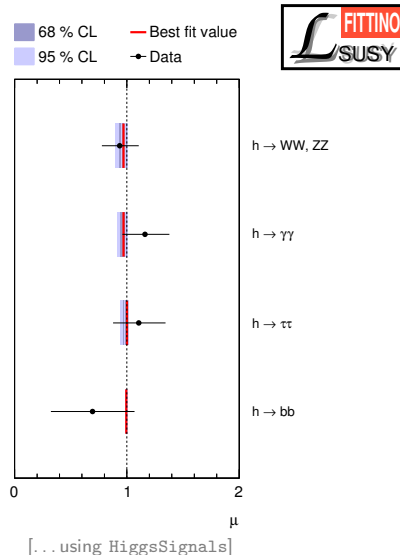
- ▶ without $(g - 2)_\mu$:
 $p = (51 \pm 3)\%$
- ▶ without Higgs properties:
 $p = (1.3 \pm 0.4)\%$
- ▶ including 10 Higgs rate and
4 Higgs mass measurements:
 $p = (4.9 \pm 0.7)\%$



How dead is the cMSSM? The impact of $(g - 2)_\mu$ and the Higgs

- ▶ without $(g - 2)_\mu$:
 $p = (51 \pm 3)\%$
- ▶ without Higgs properties:
 $p = (1.3 \pm 0.4)\%$
- ▶ including 4 Higgs rate and
1 Higgs mass measurements:
 $p = (8.3 \pm 0.8)\%$

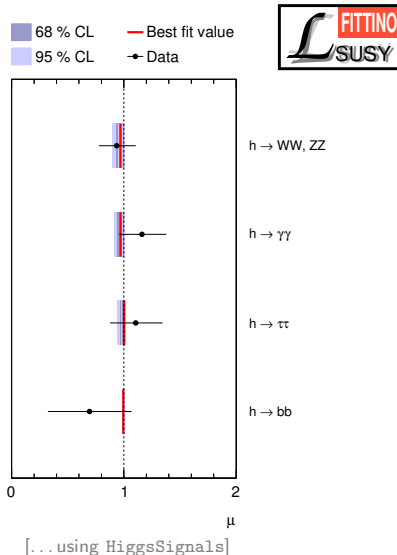
The p -value depends on the parametrization of the Higgs observables.



How dead is the cMSSM? The impact of $(g - 2)_\mu$ and the Higgs

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The cMSSM is a dull zombie...



Future global SUSY/BSM fits should

- ▶ include the Higgs observables; ✓
- ▶ calculate the p -value using toy fits; ✓
- ▶ address more general models;
- ▶ include a larger set of LHC observables.

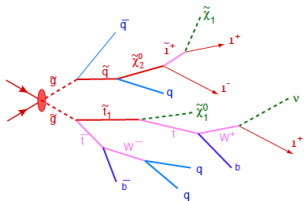
Global SUSY/BSM fits: outlook

Future global SUSY/BSM fits should

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Use **simplified models**?

Do not analyse data in terms of



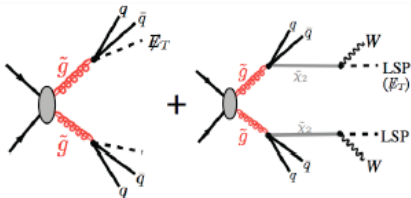
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Use [simplified models](#)?

but rather in terms of few simple topologies like

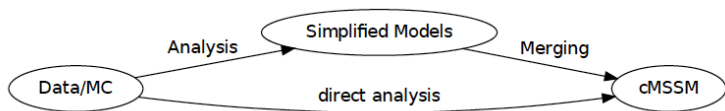


Towards model independent searches: simplified models

However...

- ▶ how do we choose the right simplified models?
- ▶ need to translate interpretation of data obtained using simplified models to interpretation of physics models
- ▶ need to assume that acceptance and detector efficiency of simplified models is similar for a specific model; how to quantify the error?
- ▶ what about backgrounds? Specific models may have additional sources of background which are not included in simplified models

The error in the translation **simplified model** \longleftrightarrow **physics model**

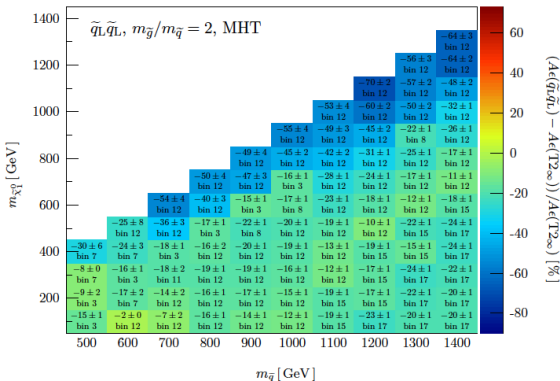


needs to be quantified

Towards model independent searches: simplified models

- Compare simplified model T2 with squark production in the MSSM
[Edelhäuser, Heisig, MK, Oymanns, Sonneveld (1410.0965)]

acceptance \times efficiency

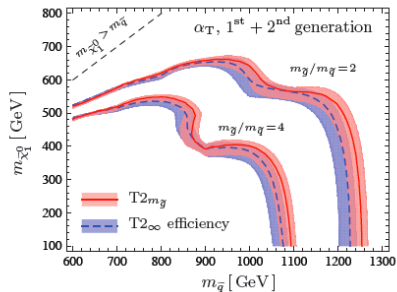
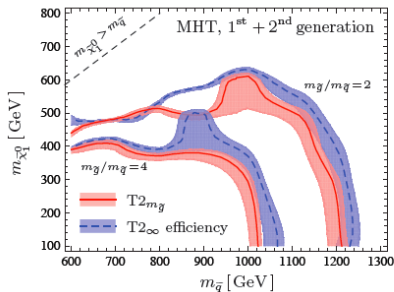


Towards model independent searches: simplified models

- Compare simplified model T2 with squark production in the MSSM

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mass limits



→ simplified models provide a reliable tool to interpret LHC searches in a more model-independent way.

Where could new physics hide?

New physics could have a

- ▶ **split spectrum** and thus reduced production cross section (e.g. "natural" SUSY with light stops)
- ▶ **different decay pattern** (e.g. compressed spectra)

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For LHC13 we need to

- ▶ explore a wider range of BSM models
- ▶ move towards more model-independent searches (effective field theories, simplified models, ...)
- ▶ increase the precision of the theoretical predictions (higher-order corrections, pdfs, ...)

However, new physics could simply be unnatural . . .

anthropic desert

Garden Eden of new physics



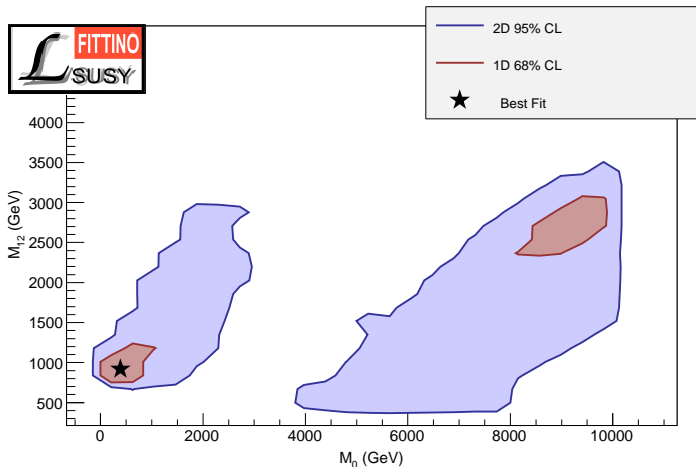
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LHC 13

Backup slides

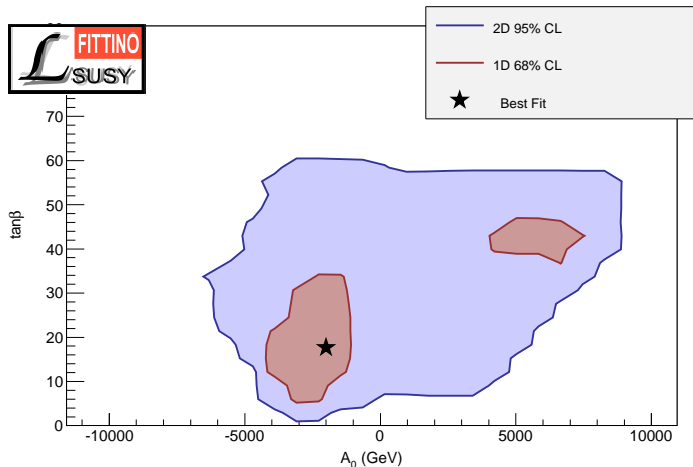
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cMSSM fit including LHC exclusions and Higgs properties:



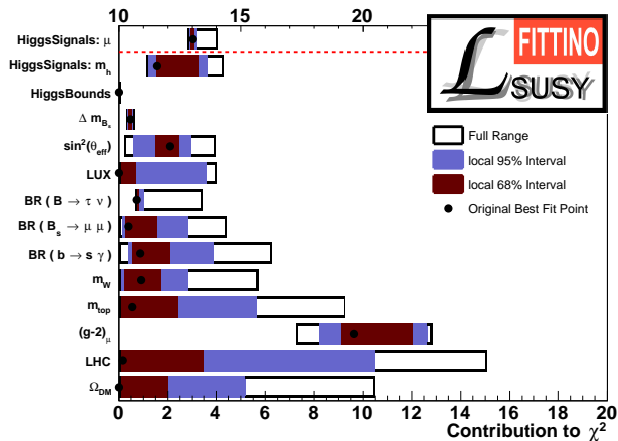
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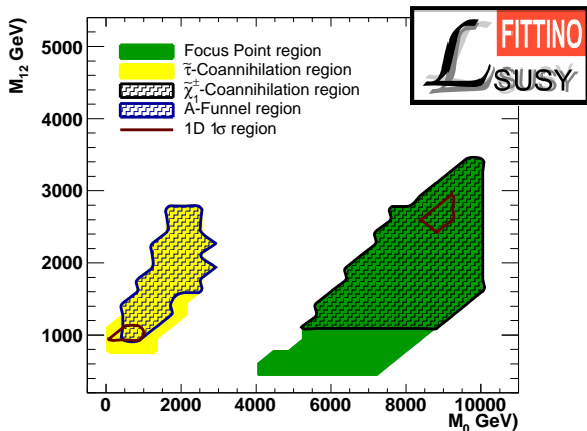
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