

A Holographic/Composite Twin Higgs

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Technion

M.G. and O. Telem, Phys.Rev.Lett. 114 (2015) 191801

C. Csaki, M.G., O. Telem and A. Weiler, in progress



Motivation

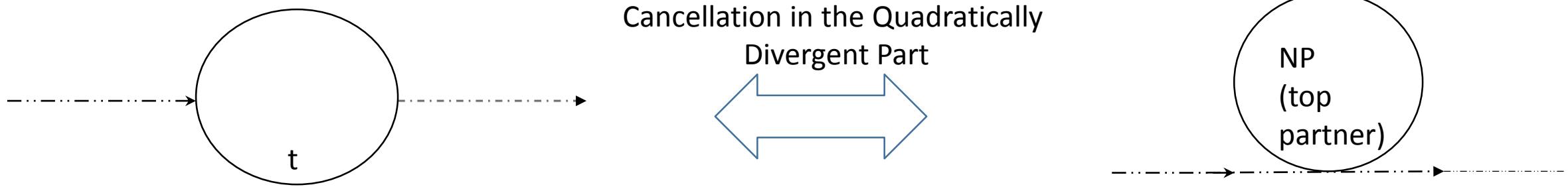
Neutral naturalness → *Natural theories with no BSM @ LHC*

The first neutrally natural model – Twin Higgs

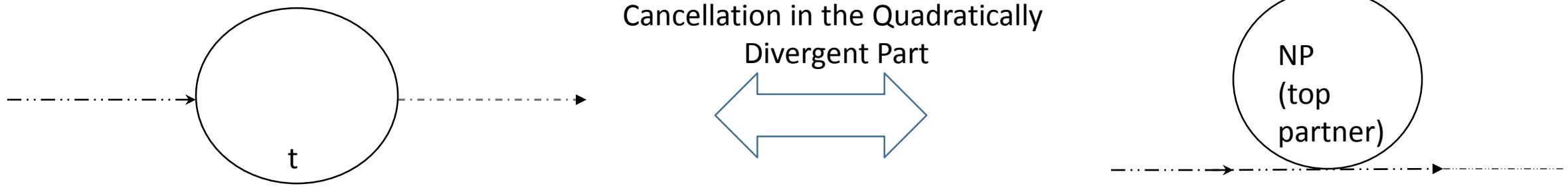
Twin Higgs needs a UV completion – *composite/AdS*

Naturalness \leftrightarrow Colored BSM

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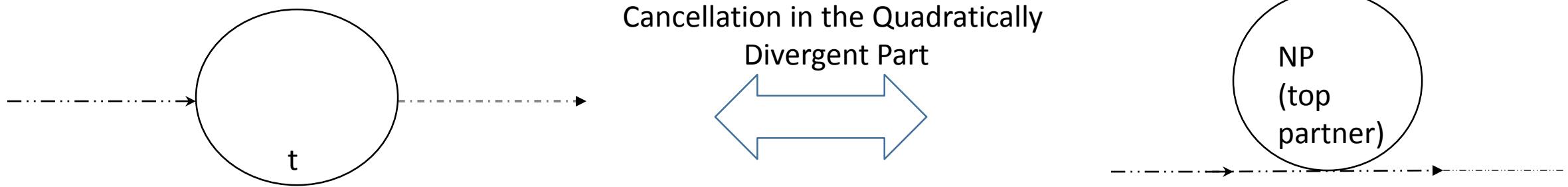


Naturalness \leftrightarrow Colored BSM



The argument:

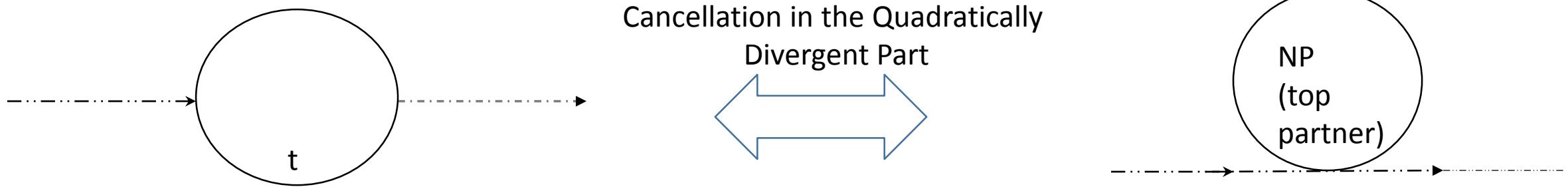
Naturalness \leftrightarrow Colored BSM



The argument:

- A symmetry is required connecting top \leftrightarrow top partners

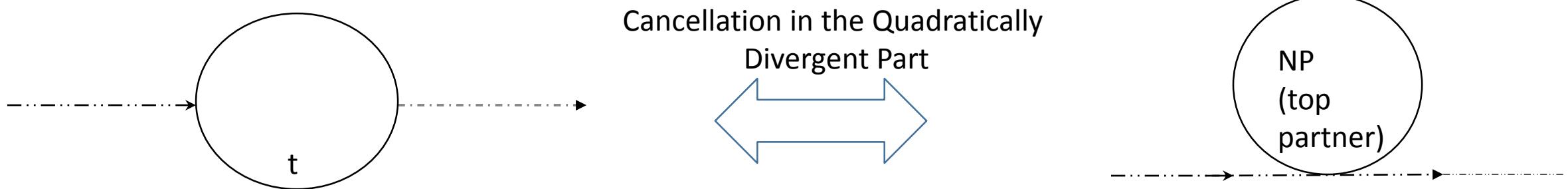
Naturalness \leftrightarrow Colored BSM



The argument:

- A symmetry is required connecting top \leftrightarrow top partners
- Naturalness requires top partners @ 1 TeV

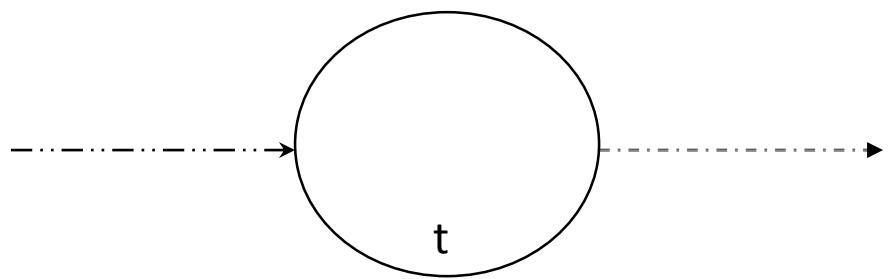
Naturalness \leftrightarrow Colored BSM



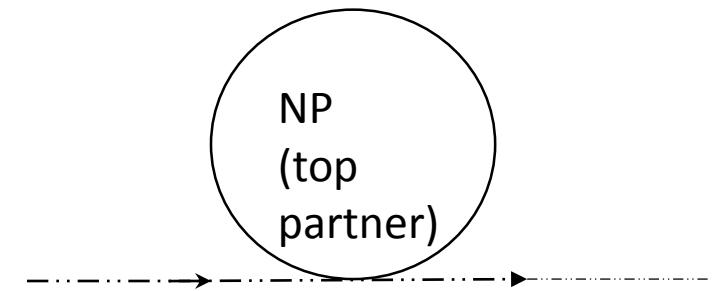
The argument:

- A symmetry is required connecting top \leftrightarrow top partners
- Naturalness requires top partners @ 1 TeV
- Colored BSM @ \sim 1 TeV

“Loophole”



Cancellation in the Quadratically
Divergent Part



top partners don't have to be colored! Just need the $N_c=3$ factor.

The Twin Higgs Model

Z. Chacko, H. S. Goh and R. Harnik, Phys. Rev. Lett. 96 (2006) 231802

Bottom-up approach: N. Craig, A. Katz, M. Strassler, R. Sundrum ,[arXiv:1501.05310](https://arxiv.org/abs/1501.05310)

A global $SU(4)$ symmetry broken by H in the fundamental: $SU(4)/SU(3)$

Gauge the group:

$$SU(2)^A \times SU(2)^B$$

SM Mirror

$$H = \begin{pmatrix} 0 \\ 0 \\ 0 \\ f \end{pmatrix}$$

$$H = \begin{pmatrix} H_A \\ H_B \end{pmatrix}$$

7 Goldstones: 6 Eaten and 1 Higgs (Pseudo-Goldstone)

Impose a Z_2 symmetry $SM \leftrightarrow Mirror$.

The Twin Higgs Model: Higgs Potential

Gauging the $SU(2) \times SU(2)$ breaks the $SU(4)$

$$\Delta V = \frac{9g_A^2\Lambda^2}{64\pi^2}H_A^\dagger H_A + \frac{9g_B^2\Lambda^2}{64\pi^2}H_B^\dagger H_B \xrightarrow{Z_2} \frac{9g^2\Lambda^2}{64\pi^2}H^\dagger H$$

SU(4) symmetric
does not produce a Goldstone mass.

Quadratically divergent terms cancel!

To have the same effect for the top loop: **double the SM symmetry**

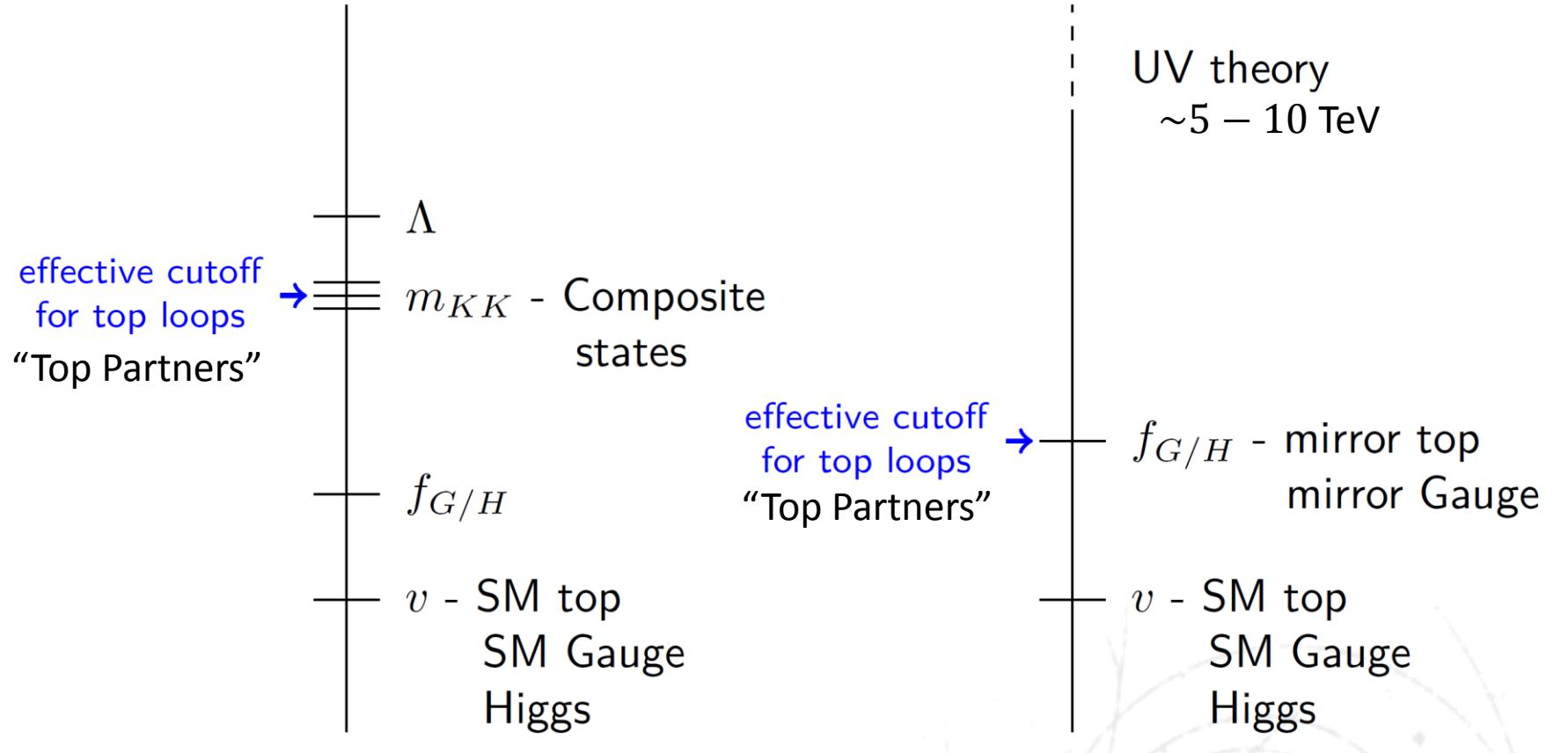
$$\begin{array}{c} (SU(3) \times SU(2) \times U(1))^A \\ \text{SM} \end{array} \times \begin{array}{c} (SU(3) \times SU(2) \times U(1))^B \\ \text{"Mirror" SM} \end{array}$$

$$H = \begin{pmatrix} 0 \\ v \\ 0 \\ f \end{pmatrix}$$

Top partners are SM singlets – "Mirror Partners"!

$$m_t^m = \frac{f}{v} m_t$$

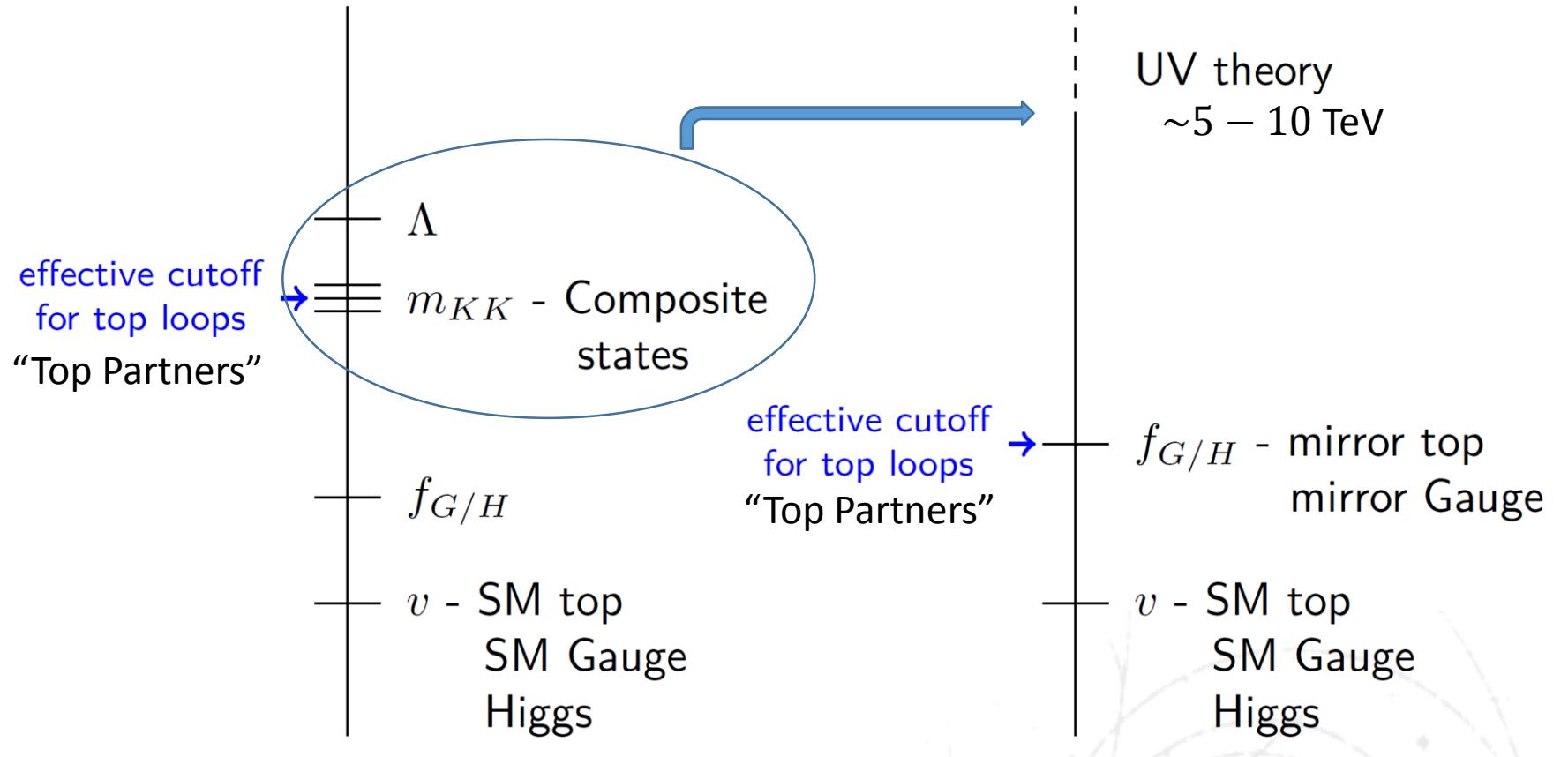
Twin Higgs and Composite Higgs



Composite Higgs
Naturalness of a PNGB Higgs
Requires light Composite states –
charged/colored naturalness

Twin Higgs
Solution to the little hierarchy
problem. “Neutral Naturalness”

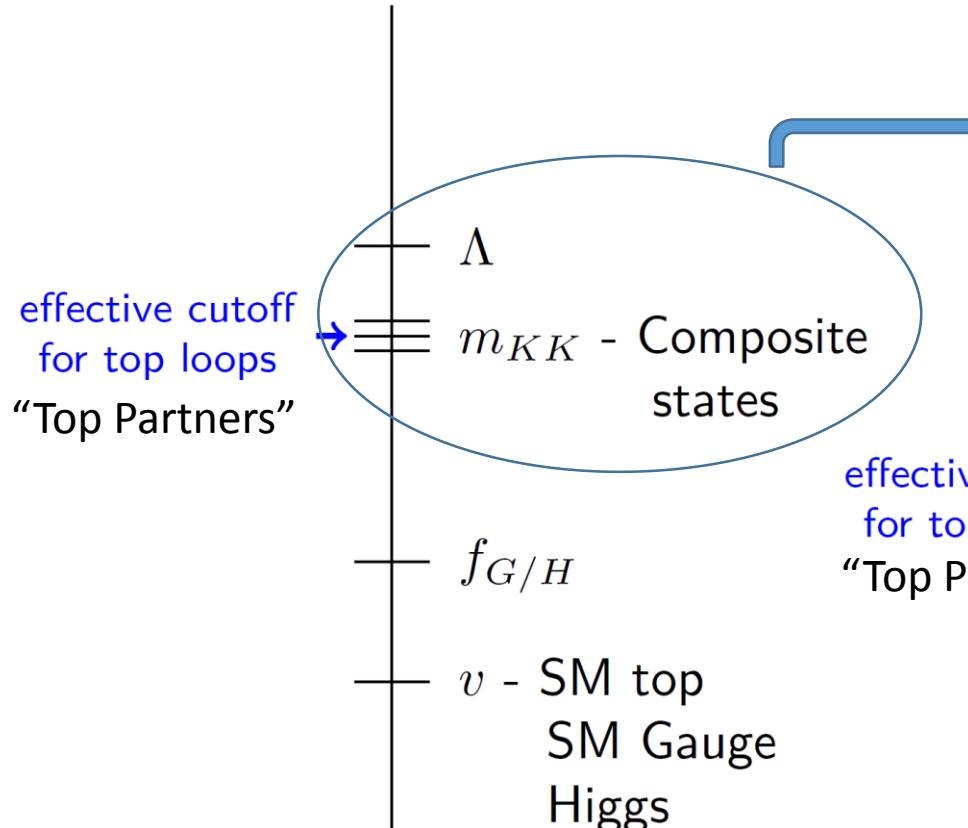
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Composite Higgs
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SUSY:

- N. Craig and K. Howe JHEP 1403 (2014) 140
A. Falkowski, S. Pokorski, M. Schmaltz, Phys.Rev. D74(2006) 035003;
S. Chang , L. J. Hall, N. Weiner Phys.Rev. D75 (2007) 035009

Orbifold:

- N. Craig, S. Knapen, P. Longhi, JHEP 1503 (2015) 106
N. Craig, S. Knapen, P. Longhi , Phys.Rev.Lett. 114 (2015) 6, 061803

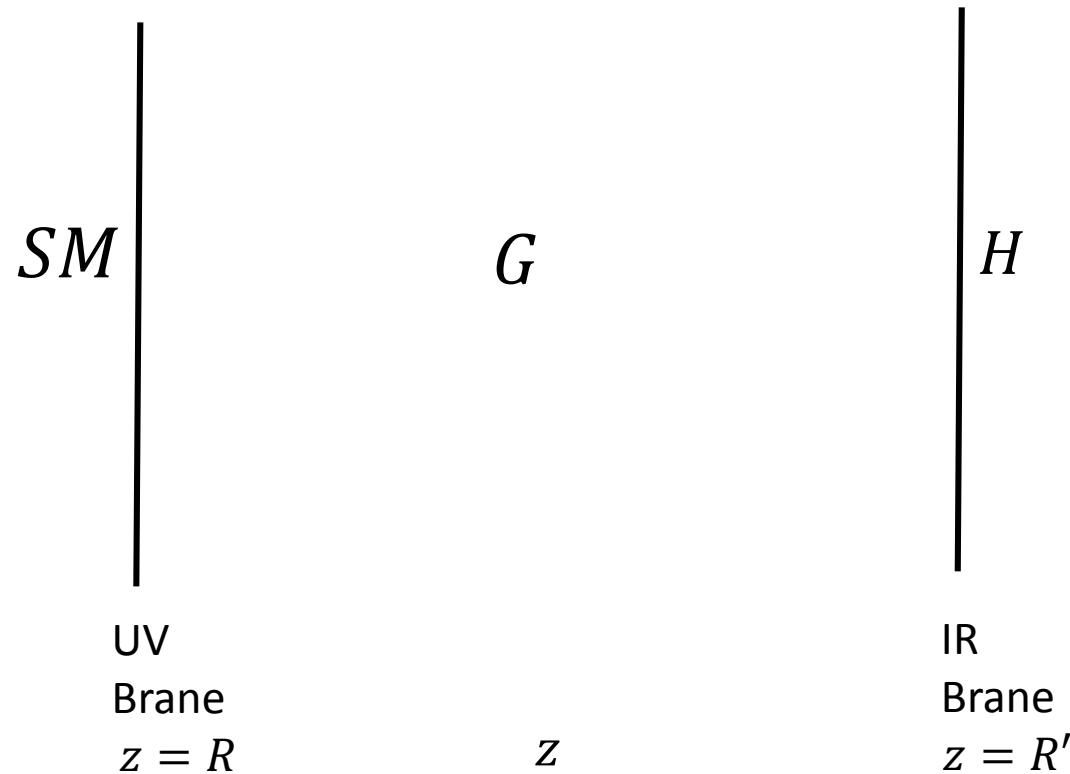
→ $f_{G/H}$ - mirror top
mirror Gauge

v - SM top
SM Gauge
Higgs

Twin Higgs
Solution to the little hierarchy
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Composite Higgs (Gauge-Higgs Unification)

- The holographic dual of composite Higgs:



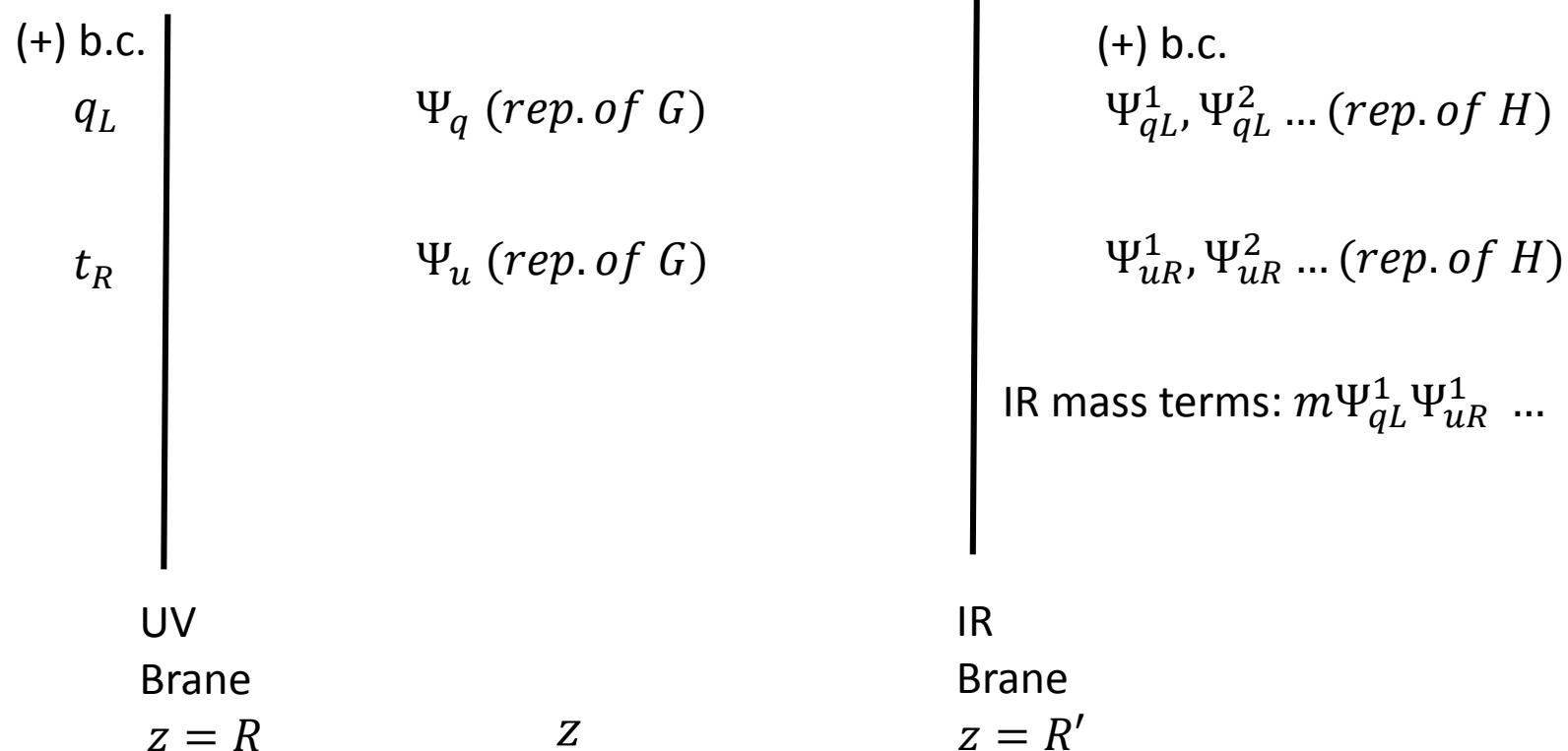
The Higgs is the fifth component of the G/H gauge fields:

$$A_5(z) = \sqrt{\frac{2}{R} \frac{z}{R'}} T_{G/H}^a h^a$$

The A_5 is a zero mode at tree level, and gets potential due to (mostly) top and SM gauge loops.

Gauge Higgs Unification

- To find the Higgs potential we need to find the masses of the KK modes.



Gauge Higgs Unification

The Gauge-Higgs vev enters the fermion EOMs:

$$\Psi_q(z, \nu) = \Omega(z, \nu)\Psi_q(z) \quad \Omega(z) = e^{ig_5 \int A_5(z)} \text{ - The Wilson line}$$

With some definitions:

$$g_* \triangleq g_5 \sqrt{R}$$

$$f \triangleq \frac{2}{g_* R'}$$

$$M_{KK} \triangleq \frac{2}{R'} = g_* f$$

$$\Omega(R') = e^{\frac{iT^a h^a}{f}\sqrt{2}} \text{ - The Goldstone matrix}$$

The Higgs Potential

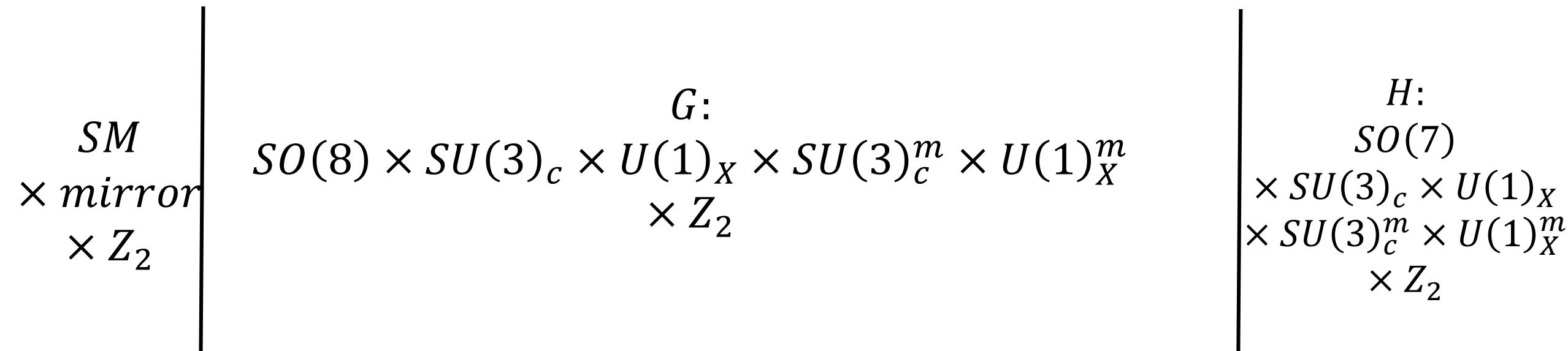
The Coleman-Weinberg potential for the Higgs is calculated using:

$$V(h) = \frac{N}{(4\pi)^2} \int dpp^3 \log(\rho[-p^2])$$

$\rho(p^2)$ is the spectral function –

$\rho(m_n^2) = 0$ for any KK state in the presence of the EW vacuum.

The Holographic Twin Higgs



$$SO(8) \rightarrow SU(2)_L \times SU(2)_R \times SU(2)_L^m \times SU(2)_R^m \quad Y = T_R^3 + X \quad Y^m = T_R^{3m} + X^m$$

Possibly: $SU(3)_c \times U(1)_x \times SU(3)_c^m \times U(1)_x^m \times Z_2 \in SU(7)$

The Top Quark

	$SO(8)$	$SU(3)_c$	$U(1)_X$
Ψ_q	8	3	2/3
Ψ_t	1	3	2/3

bulk Z_2 , 7 of $SU(7)$

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UV b.c. (+):

	$SU(2)_L$	$SU(3)_c$	$U(1)_Y$
Q_L	2	3	1/6
t_R	1	3	2/3

UV Z_2

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Ψ_{qL}^7, Ψ_{qL}^1	7,1	3	2/3
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IR Z_2 , 7 of $SU(7)$

	$SO(7)$	$SU(3)_c^m$	$U(1)_X^m$
$\Psi_{qL}^{7m}, \Psi_{qL}^{1m}$	7,1	3	2/3
Ψ_{tR}^m	1	3	2/3

IR masses: $m_t^1 \Psi_{qL}^1 \Psi_{tR}$

IR masses: $m_t^1 \Psi_{qL}^{1m} \Psi_{tR}^m$

The spectral function of Composite-Twin Higgs

- The spectral functions of the top and mirror top:

$$\rho_t(p^2) = 1 + f_t(p^2) \sin^2 \left(\frac{h}{f} \right)$$

$$\rho_{tm}(p^2) = 1 + f_t(p^2) \cos^2 \left(\frac{h}{f} \right)$$

- The Higgs potential

$$V_{eff}(h) = \frac{-4N_c}{(4\pi)^2} \int_0^\infty dp p^3 \log(\rho_t[-p^2] \rho_{tm}[-p^2])$$

The top mass

- At low energies

$$\rho_t(p) \approx 1 - m_t^2/p^2 \quad \rho_{tm}(p) \approx 1 - m_{tm}^2/p^2$$

- The top mass can be calculated

$$m_t = \frac{\frac{g_* v}{2\sqrt{2}} \tilde{m}_t f_q f_{-u}}{\sqrt{1 + f_u^2 f_{-q}^{-2} \tilde{m}_t^2}}$$
$$f_c = \sqrt{\frac{1-2c}{1-\left(\frac{R'}{R}\right)^{2c-1}}}$$

- Using the known $m_t(3 \text{ TeV})$ we can find $g_*(c_q, c_u, \tilde{m}_t)$

The Higgs Potential

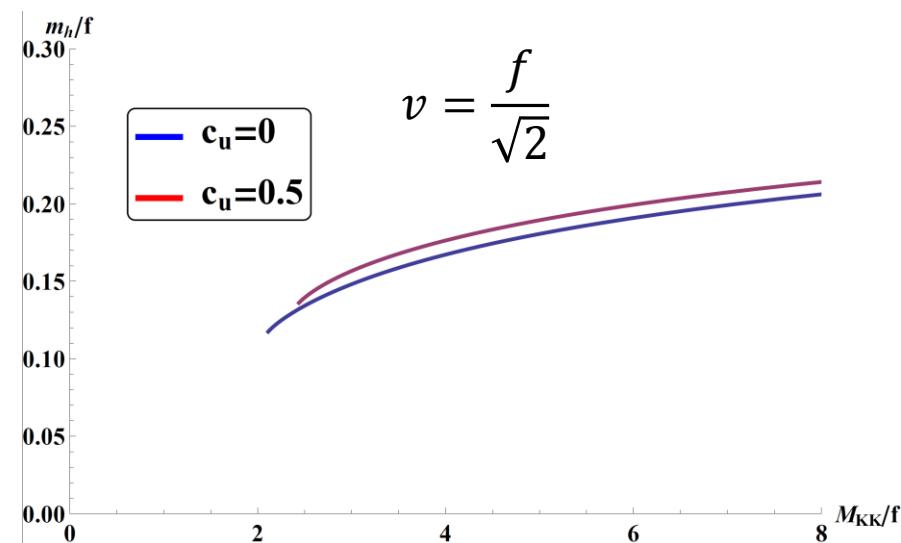
The Higgs potential:

$$V(h) \approx -\alpha_2 \sin^2 \frac{h}{f} + \frac{\alpha}{2} \sin^4 \frac{h}{f} - \alpha_2 \cos^2 \frac{h}{f} + \frac{\alpha}{2} \cos^4 \frac{h}{f} = \alpha \sin^2 \frac{h}{f} \cos^2 \frac{h}{f} + const$$

$\underbrace{\hspace{10em}}$ Top+Gauge $\underbrace{\hspace{10em}}$ Mirror top+gauge

$$\alpha_2 \sim \frac{3}{32\pi^2} y_t^2 f^2 m_{KK}^2$$

$$\alpha \sim \frac{3}{64\pi^2} y_t^4 f^4 \log \frac{2M_{KK}^2}{y_t^2 f^2}$$



The Higgs Potential

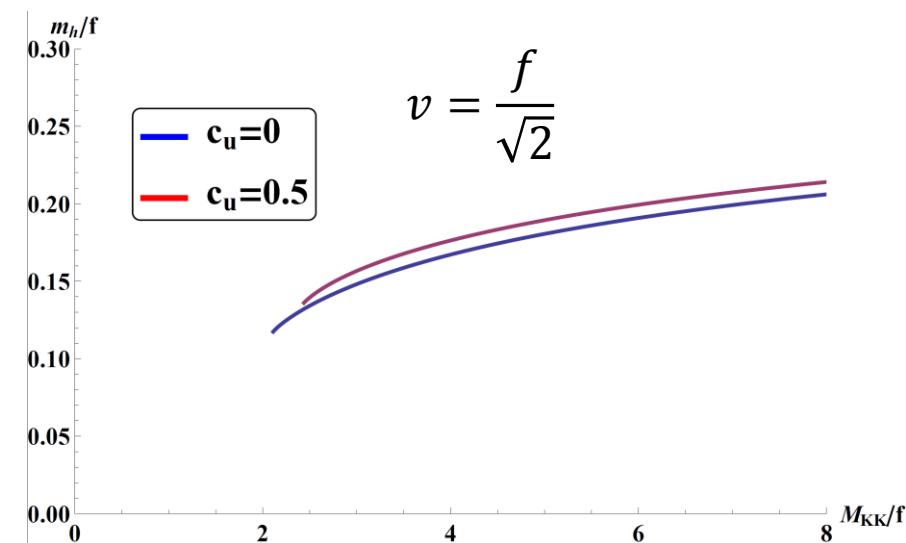
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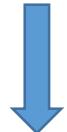
$$\alpha \sim \frac{3}{64\pi^2} y_t^4 f^4 \log \frac{2M_{KK}^2}{y_t^2 f^2}$$



The Higgs Potential

Suppose we have added a term:

$$V(h) = -\alpha \sin^2 \frac{h}{f} \cos^2 \frac{h}{f} + \beta \sin^2 \frac{h}{f}$$



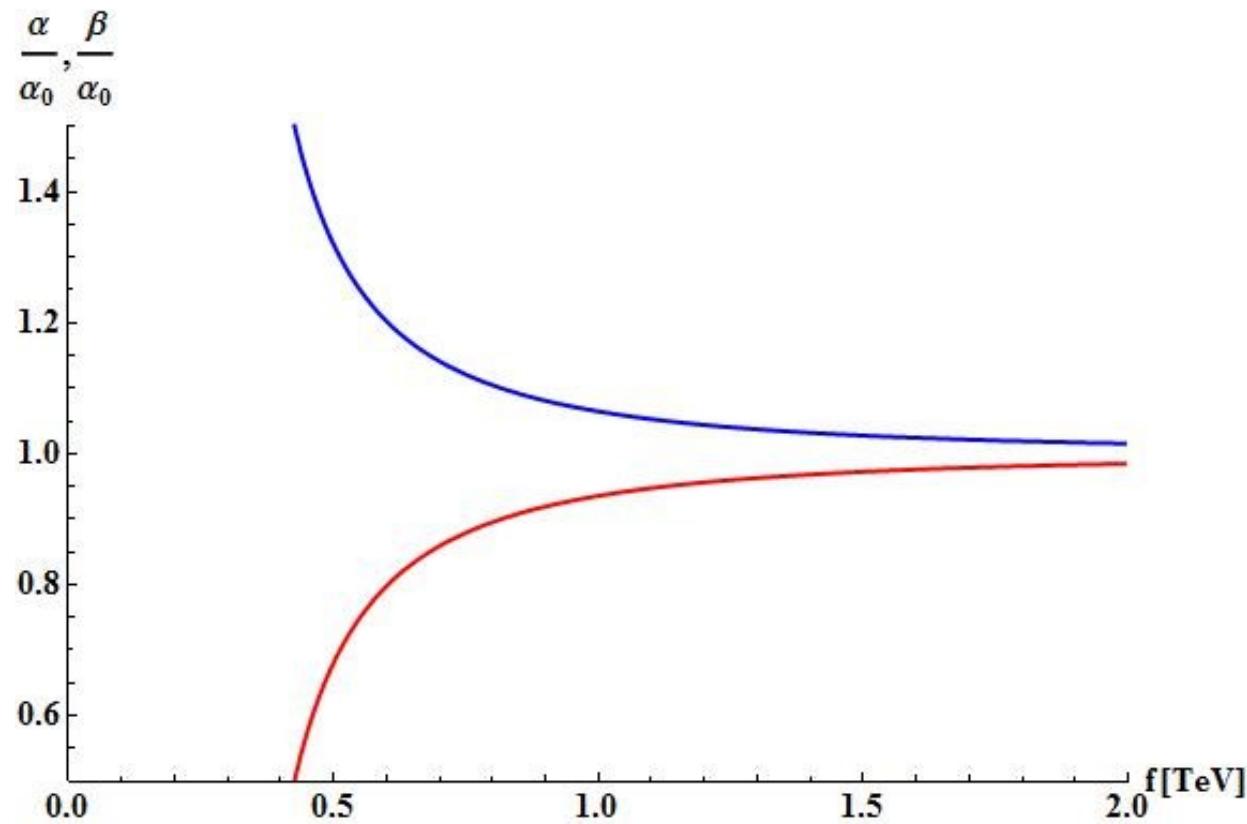
Higgs mass and vev

$$\frac{\alpha}{\alpha_0} = \frac{1}{1 - \epsilon^2} \quad \frac{\beta}{\alpha_0} = \frac{1 - 2\epsilon^2}{1 - \epsilon^2}$$

$$\alpha_0 = \frac{f^4 m_h^2}{8v^2} \quad \epsilon = v/f$$

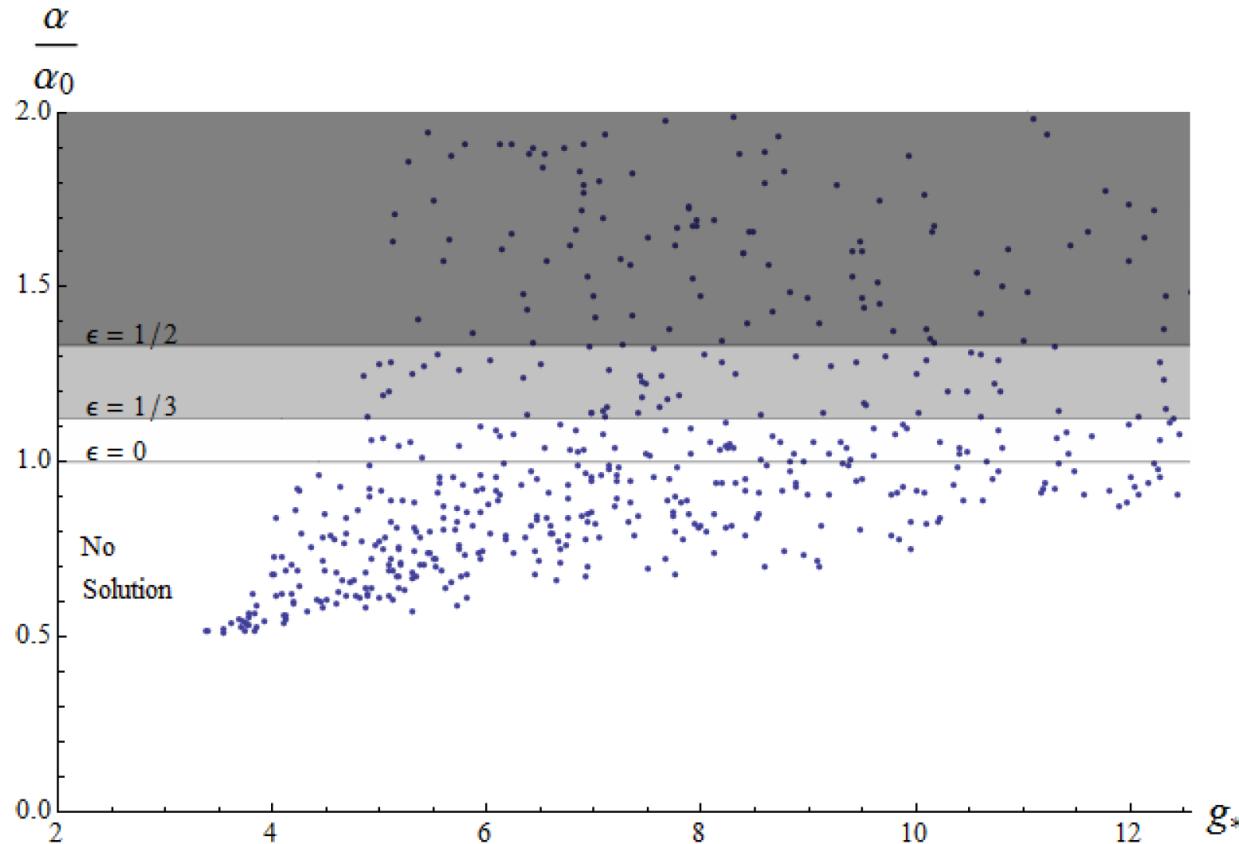
- The tuning is:

$$\Delta = \max \left| \frac{d \log v}{d \log p_i} \right| = \frac{1}{2} \left(\frac{f^2}{2v^2} - 1 \right) \max \left| \frac{d \log \alpha, \beta}{d \log p_i} \right|$$



The Higgs Potential

- The Z_2 conserving contribution - $\alpha(c_q, c_u, m)$



$$\Delta \sim \frac{f^2}{v^2}$$

Holographic Twin Higgs

Naturalness without KK – tops

- The KK tops don't enter the tuning.
- The tuning scales as f^2/v^2 – Higgs data.
- M_{KK} can be arbitrarily high, but unitarity requires $M_{KK} < 4\pi f$

Need Z_2 breaking

- Higgs potential – $\beta \sin^2 \frac{h}{f}$
- Dark radiation – Mirror neutrinos and Mirror photon.

Pheno

EW precision

Tree Level: $M_{KK} > 3 \text{ TeV}$

Higgs Loops: may potentially be dangerous

Vector-like Quarks/Resonances

LHC reach: 1.5 TeV for Kktops
 ~4 TeV for KKGlue

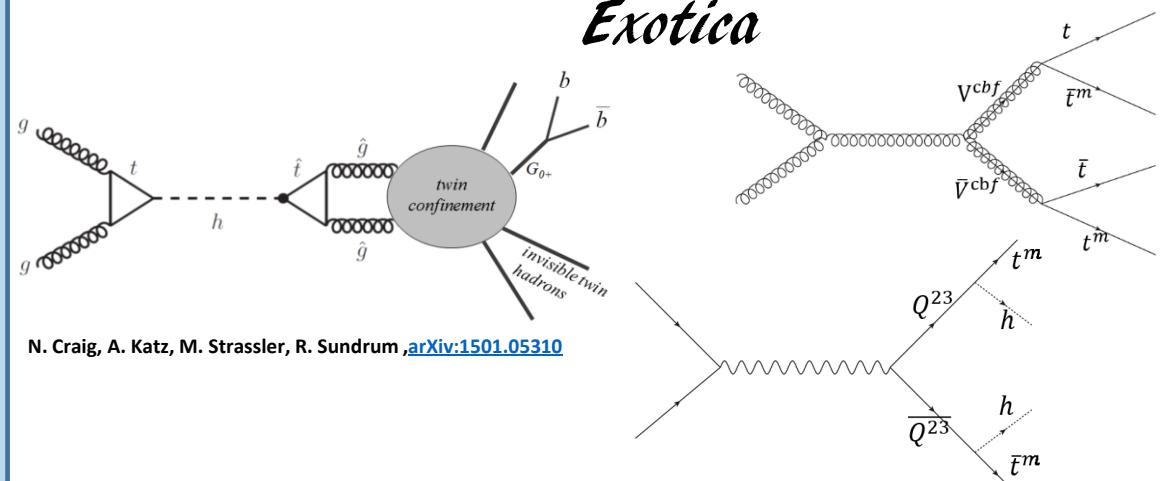
Accessible in future hadronic colliders

Higgs precision

PNGB - all couplings $\left(1 - \frac{v^2}{f^2}\right)$

Invisible Decays $Br(h \rightarrow b^m b^m) \approx \frac{v^2}{f^2} Br(h \rightarrow bb)$

Exotica



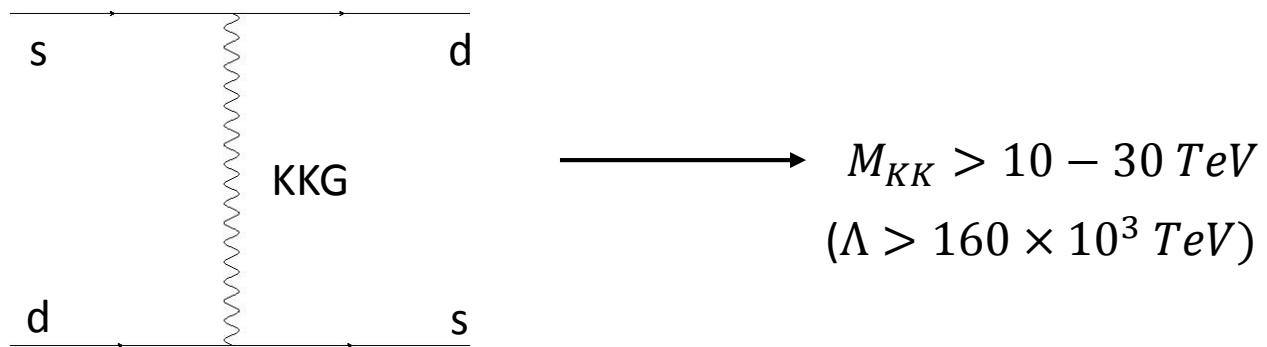
Flavor – Work in Progress

In composite Higgs flavor violation is suppressed by RS-GIM:

- Flavor violation is due to KK modes - occurs near the IR brane.
- The wavefunction of the light quarks on the IR brane is suppressed:

$$f_q \sim f_{-u} \sim \sqrt{\frac{m_q}{g_* v}}$$

Anarchic flavor in composite Higgs is not enough



See e.g. C. Csaki et. al. arXiv:0804.1954

Flavor in Composite Twin Higgs

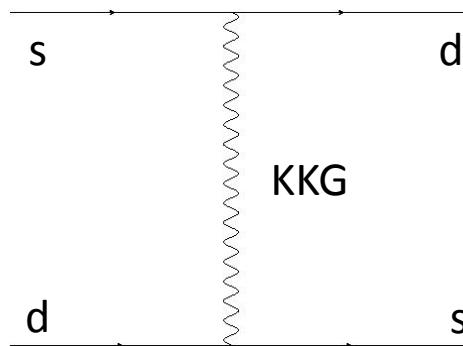
Contrary to CH – M_{KK} is bound only by pertubativity: $M_{KK} < 4\pi f$

Is anarchic flavor enough in the CTH? What is the limit on f ?

Flavor in Composite Twin Higgs

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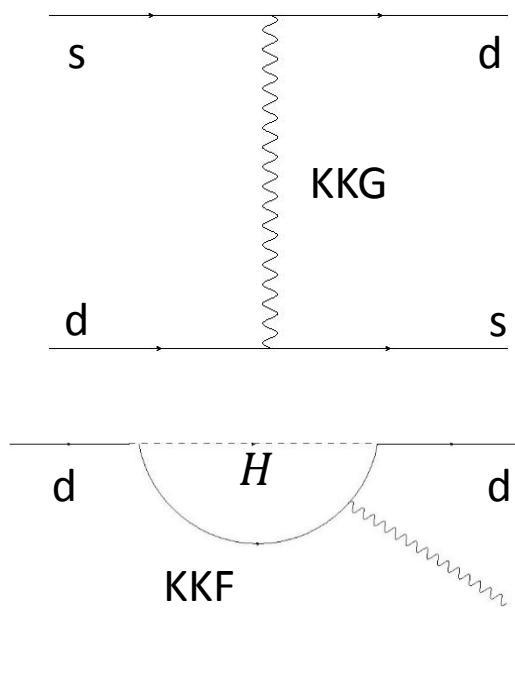
$$C_k^4 = \frac{1}{(1.6 \times 10^5 \text{ TeV})^2} \left(\frac{106 \text{ TeV}}{g_*^2 f \tilde{m}_d} \right)^2$$

$$g_*^2 f \tilde{m}_d > 106 \text{ TeV}$$

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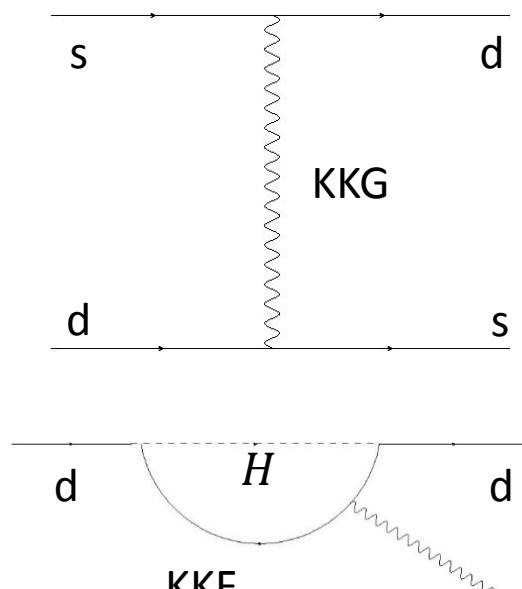
$$\frac{c}{16\pi^2 f^2} m_d \bar{d}_L \sigma^{\mu\nu} e F_{\mu\nu} d_R \quad c \sim \frac{1}{g_*^2} (Y^2) = \frac{\tilde{m}_d^2}{4}$$

$$\frac{2f}{\tilde{m}_d} > 7.29 \text{ TeV}$$

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$$\frac{2f}{\tilde{m}_d} > 7.29 \text{ TeV}$$

$$g_* f > 19.6 \text{ TeV}$$

Flavor in Composite Twin Higgs

In our estimations the combination of the two most stringent bounds:

$$g_* f > 19.6 \text{ TeV}$$

At the limit of perturbative control $g_* = 4\pi \rightarrow f > 1.5 \text{ TeV}$.

The tuning in the Higgs potential is at the percent level.

Z_2 breaking

Bulk \mathbb{Z}_2 :

Bulk masses: $c_q \neq c_q^m \rightarrow y_q \neq y_q^m$

Gauge couplings: $g_5^{SU(3)_c/U(1)_X} \neq g_5^{SU(3)_c^m/U(1)_X^m}$

IR \mathbb{Z}_2 :

IR masses $m_{IR} \neq m_{IR}^m \rightarrow y_q \neq y_q^m$

IR b.c.: e.g. break $U(1)_X^m$ - massive mirror photon

UV \mathbb{Z}_2 :

Neumann b.c.: eliminate mirror zero modes

UV Kinetic terms $\rightarrow y_q \neq y_q^m$

Z_2 breaking

Light mirror fermions:

- Eliminate using UV \cancel{Z}_2
- Arbitrary masses – e.g. MeV-GeV mirror neutrinos using bulk \cancel{Z}_2 .

Massive mirror photon:

- UV brane – eliminated (= not gauged)
- IR brane – EW scale mass.

Higgs potential

Generate the right term using the hypercharge: $g' \neq g'_m$

$$V(h) = \alpha_2 \sin^2 \frac{h}{f} + \alpha_2^m \cos^2 \frac{h}{f} = \beta \sin^2 \frac{h}{f} + \text{const}$$

$$\beta = \alpha_2^m - \alpha_2$$

Z_2 breaking – Higgs Potential

- Hypercharge -

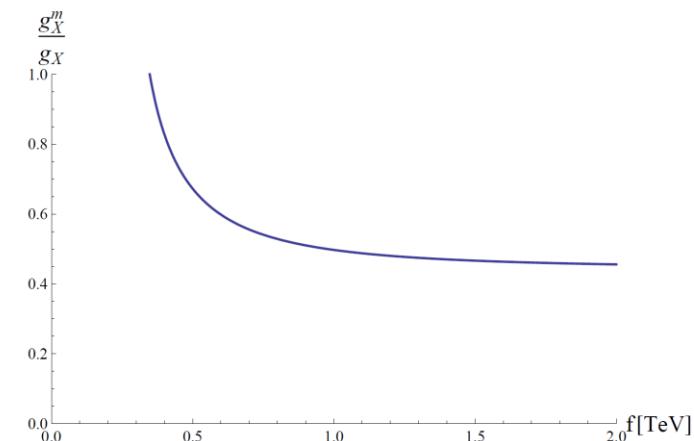
$$\frac{1}{g'^2} = \log \frac{R'}{R} \left(\frac{1}{g_*^2} + \frac{1}{g_{X*}^2} \right) \approx \frac{1}{g_{X*}^2} \log \frac{R'}{R}$$

- Detune the $U(1)_X$ gauge coupling in the bulk

$$g_{X*} > g_{X*}^m$$



$$\beta \approx \frac{3}{128\pi^2} (g'^2 - g_m'^2) g_*^2 f^4 \approx \frac{3}{128\pi^2} \frac{(g_{X*}^2 - g_{X*}^{m2})}{\log \frac{R'}{R}} g_*^2 f^4$$



Summary

- Twin Higgs – SM singlet top partners.
- Needs “UV completion” - Holographic Twin Higgs:
 - Only log dependence on M_{KK} in the Higgs potential.
 - Tuning scales as $\frac{f^2}{v^2}$. $M_{KK} \lesssim 4\pi f$ by unitarity.
 - Needs Z_2 breaking.
- Flavor:
 - Anarchic flavor requires $M_{KK} \gtrsim 20 \text{ TeV}$.
 - Leads to percent level tuning.
- Pheno:
 - Higgs couplings, Invisible width.
 - Colored states and exotics at future hadronic colliders

Thank You!



The Spectral Function in CTH

$$f_t = \frac{\frac{1}{2}C_{-1}\tilde{m}_u^2}{(C_{-8}S_1 + C_{-1}S_8\tilde{m}_u^2) S_{-8}}$$

$$C_{\pm i} \equiv C_{\pm c_i}(R', p), \quad S_{\pm i} \equiv S_{\pm c_i}(R', p)$$

$$(kz)^{c+2}C_c(z, p) = \frac{\pi p}{2k}(kz)^{\frac{5}{2}} \left[J_{c+\frac{1}{2}}\left(\frac{p}{k}\right) Y_{c-\frac{1}{2}}(zp) - Y_{c+\frac{1}{2}}\left(\frac{p}{k}\right) J_{c-\frac{1}{2}}(zp) \right]$$

$$(kz)^{c+2}S_c(z, p) = \frac{\pi p}{2k}(kz)^{\frac{5}{2}} \left[J_{\frac{1}{2}-c}\left(\frac{p}{k}\right) Y_{\frac{1}{2}-c}(zp) - Y_{\frac{1}{2}-c}\left(\frac{p}{k}\right) J_{\frac{1}{2}-c}(zp) \right]$$