

## Isotherms of liquid-gas phase transition

$$P = \frac{RT}{v-b} - \frac{a}{v^2}$$

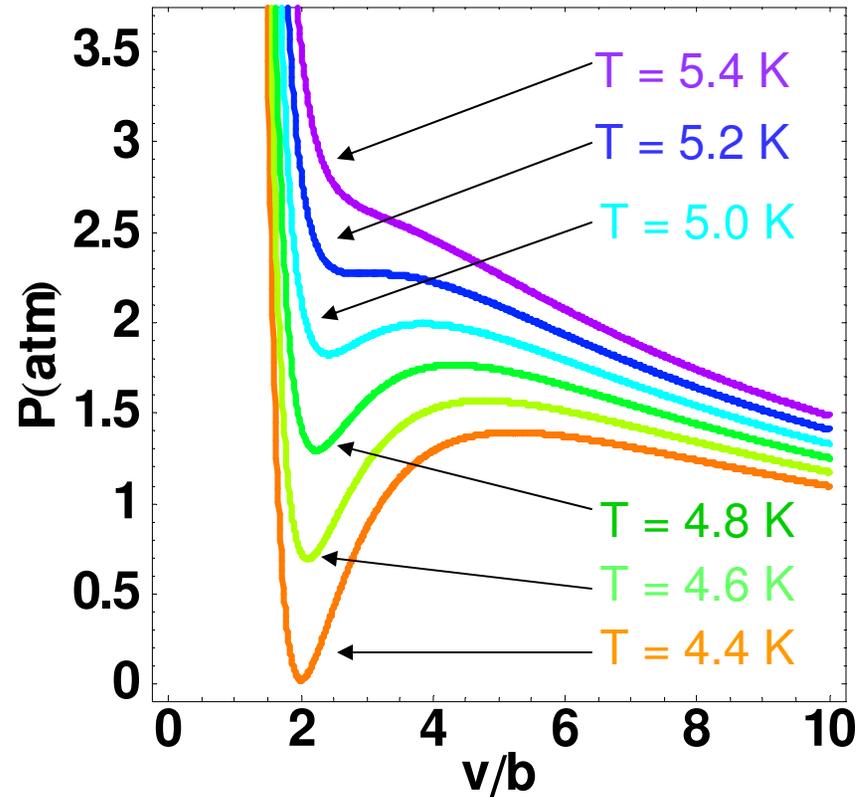
Let us plot  $P$  vs  $v$  at fixed  $T$  for vdW fluid:

For concreteness, we will use a vdW fluid approximating He.

Note that the isotherms with  $T < 5$  K do not satisfy a stability criterion:

$$\kappa_T > 0$$

$$\kappa_T = -\frac{1}{v} \left( \frac{\partial v}{\partial P} \right)_T > 0 \quad \Rightarrow \quad \left( \frac{\partial P}{\partial v} \right)_T < 0$$

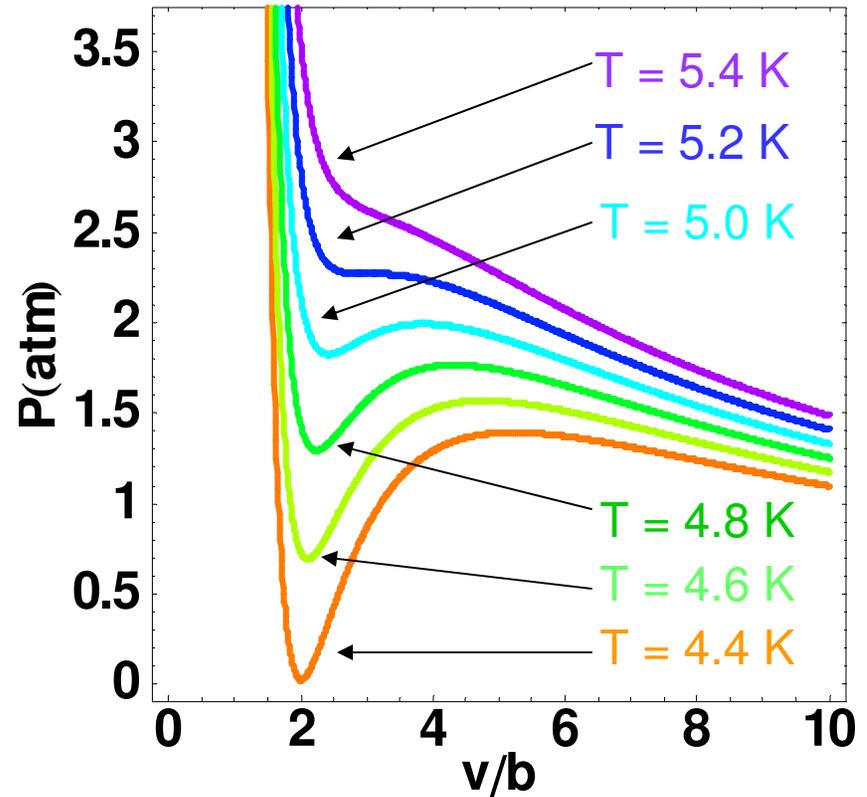


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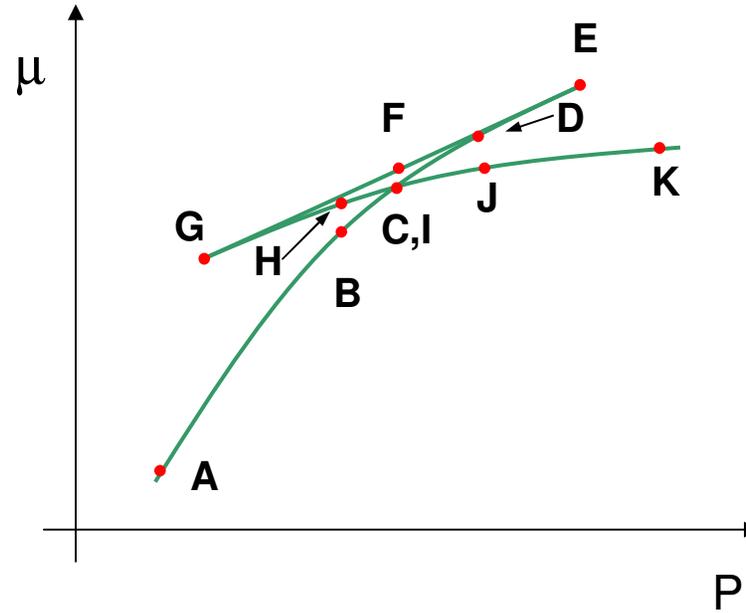
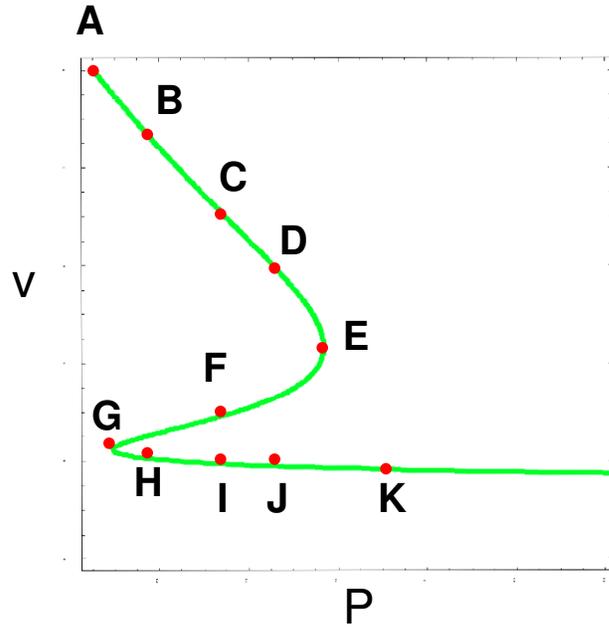
The **unstable portion** of the isotherm is unphysical and thus the **isotherm has to be modified** to describe the equilibrium state of the system.

Let us calculate the molar Gibbs potential (chemical potential) using the unstable isotherm and the Gibbs-Duhem relation:

$$d\mu = -sdT + vdP$$

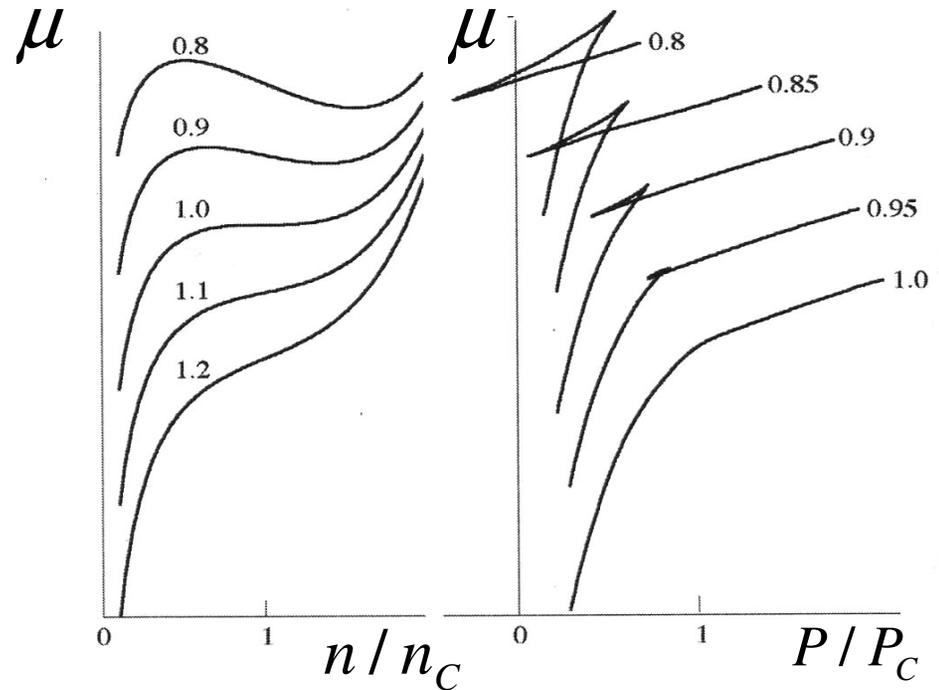
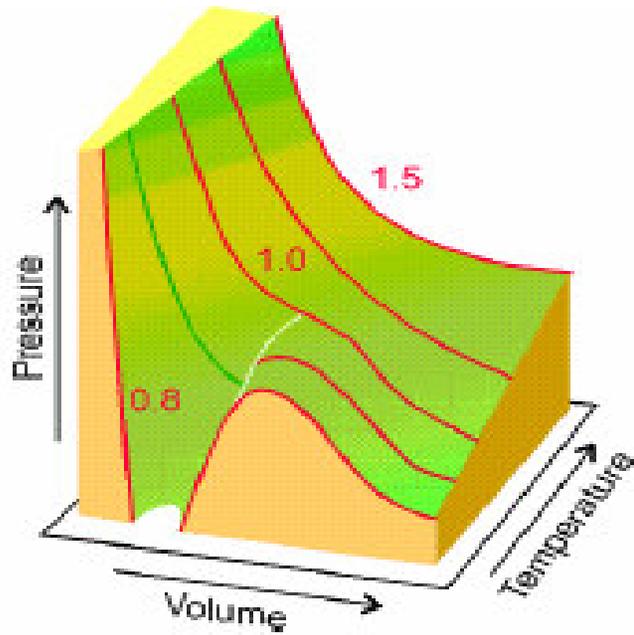


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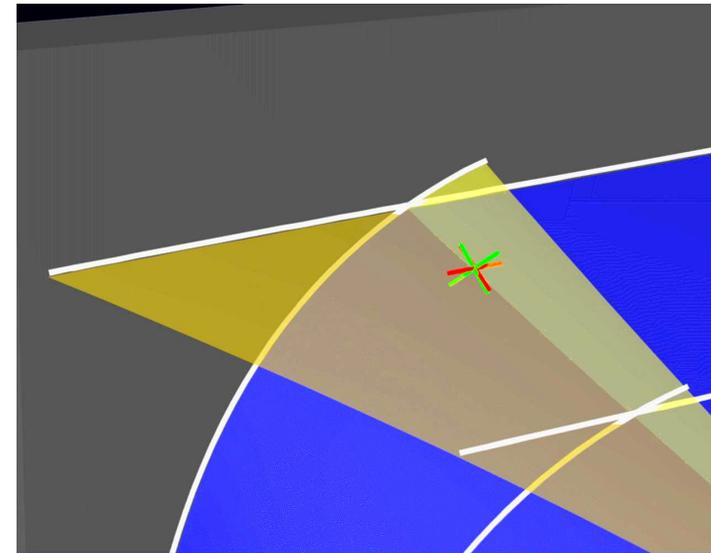
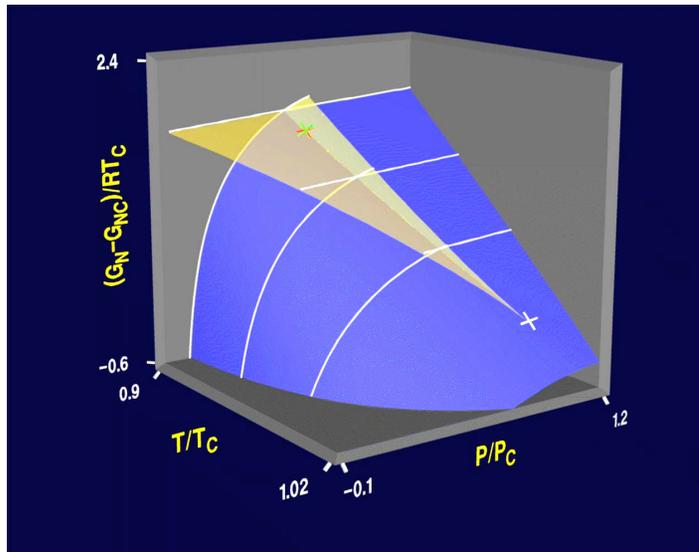


Since the Gibbs energy should be at minimum in equilibrium, all equilibrium states should lie on the curve **A-B-(C,I)-J-K**. Also note that transition between the points **C** and **I** should take place **at constant pressure and chemical potential!**

# Isotherms of liquid-gas phase transition



## 3D view of the chemical potential as a function of $P$ and $T$



Plot of the molar Gibbs potential (the chemical potential) as a function of pressure and temperature (normalized to the pressure and temperature at the critical point marked by white cross).

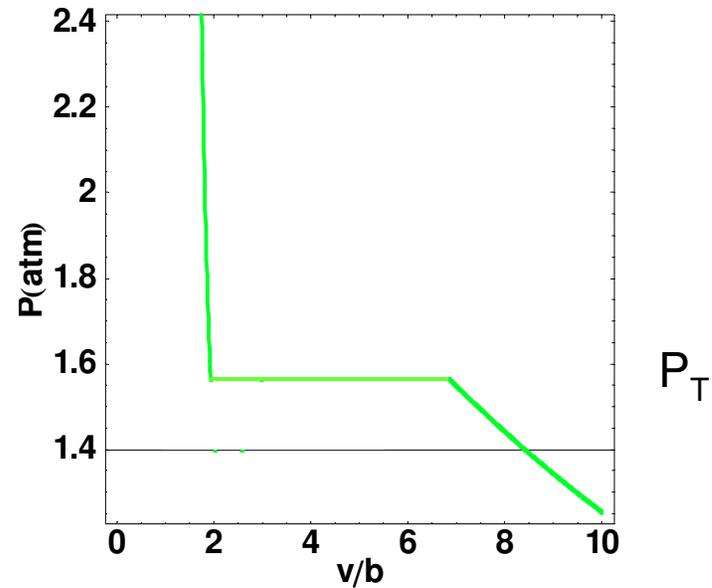
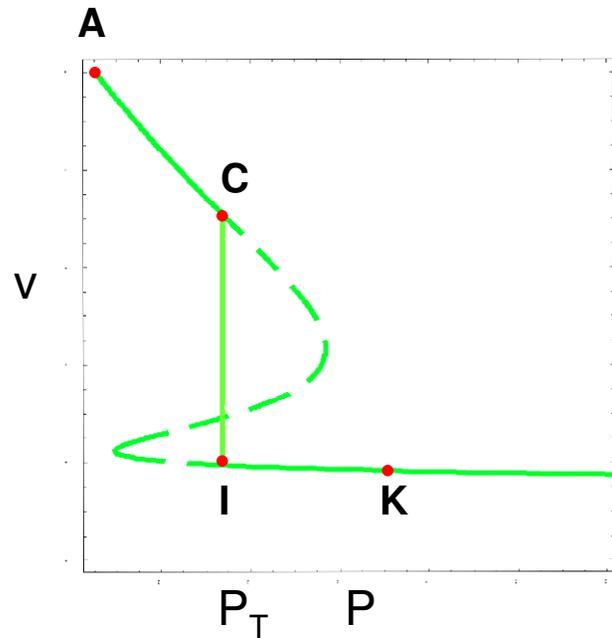
Blowup of a region of the plot on the left. The chemical potential surface intersects itself along the coexistence curve. Blue surfaces are equilibrium liquid and gas phases. Yellow surfaces are meta-stable states.

Source: [www.public.iastate.edu/~jolls/dedication.html](http://www.public.iastate.edu/~jolls/dedication.html)

Notes

## *Isotherms of liquid-gas phase transition*

Since transition between points **C** and **I** has to take place at a constant pressure, it is represented by a **straight line of  $P=\text{const}$**  on the  **$P$ - $v$**  diagram.



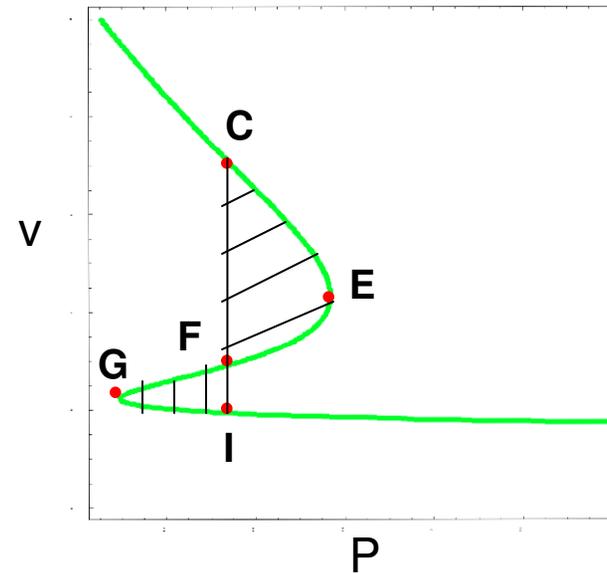
Stable isotherm for a vdW fluid

## Isotherms of liquid-gas phase transition

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The value of the pressure  $P_T$  at which the transition from **C** to **I** takes place can be determined from the condition  $\mu_C = \mu_I$ .

$$\mu_I - \mu_C = \int_C^I v(P) dP = 0$$



# Isotherms of liquid-gas phase transition

$$\mu_I - \mu_C = \int_C^I v(P) dP = 0$$



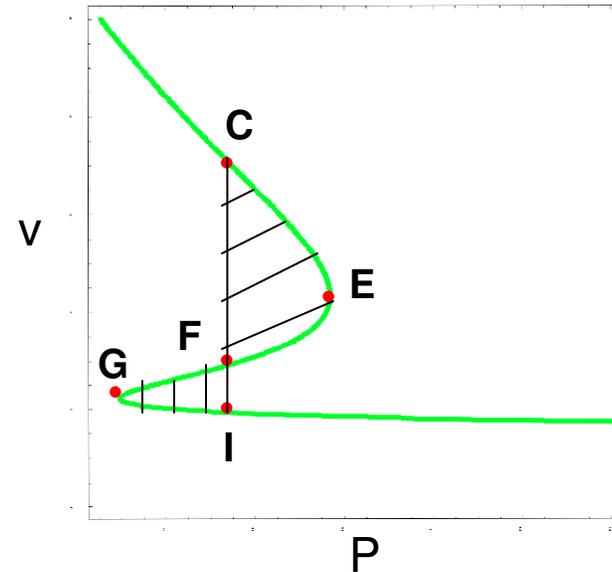
$$\mu_I - \mu_C = \int_C^E v dP + \int_E^F v dP + \int_F^G v dP + \int_G^I v dP = 0$$



$$\int_C^E v dP - \int_F^E v dP - \int_G^F v dP + \int_G^I v dP = 0$$



$$\int_C^E v dP - \int_F^E v dP = \int_G^F v dP - \int_G^I v dP$$



This condition means that the area enclosed by **CEF** should be equal to the area enclosed by **FGI**.

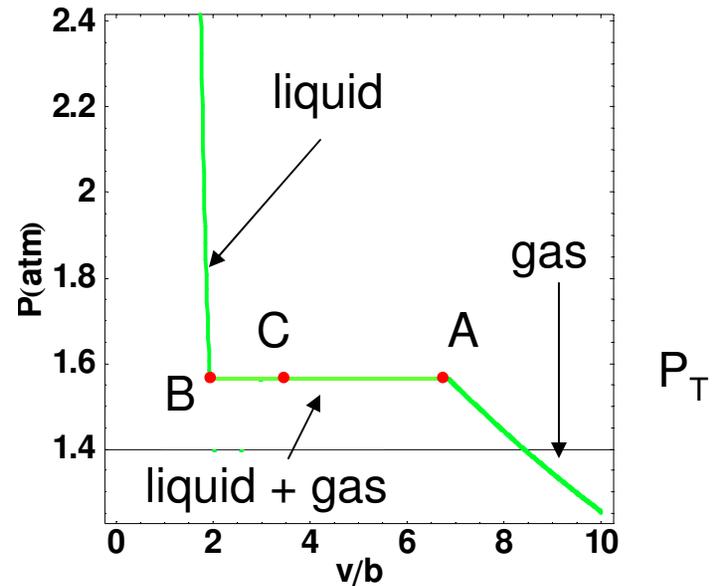
## *Isotherms of liquid-gas phase transition*

From the stable isotherm shape it is clear that there is a continuous change of molar volume across the phase transition. Molar entropy and internal energy also change across the phase transition.

$$v = x_L v_L + x_G v_G \quad x_L + x_G = 1$$

where  $v_L$  and  $v_G$  are molar volumes of liquid and gas and  $x_L$  and  $x_G$  are molar fractions of liquid and gas in the liquid/gas mixture. Solving for  $x_L$ :

$$x_L = \frac{v_G - v}{v_G - v_L} \quad x_G = \frac{v - v_L}{v_G - v_L} \quad x_L = \frac{CA}{BA} \quad x_G = \frac{BC}{BA} \quad \text{- "lever rule"}$$

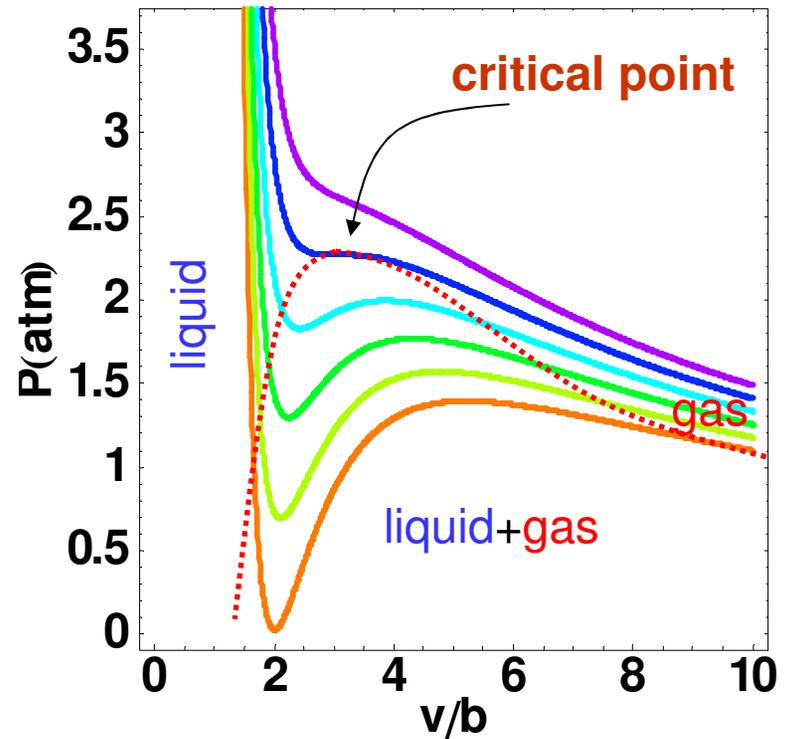


## *Isotherms of liquid-gas phase transition*

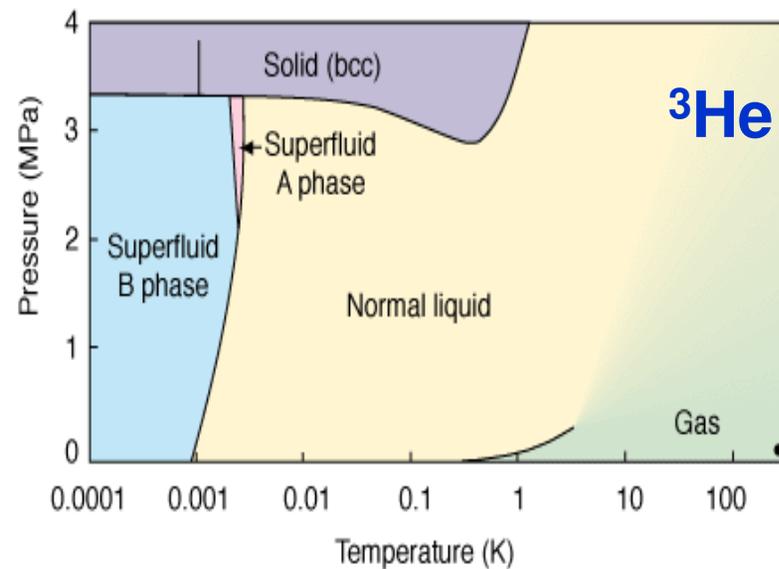
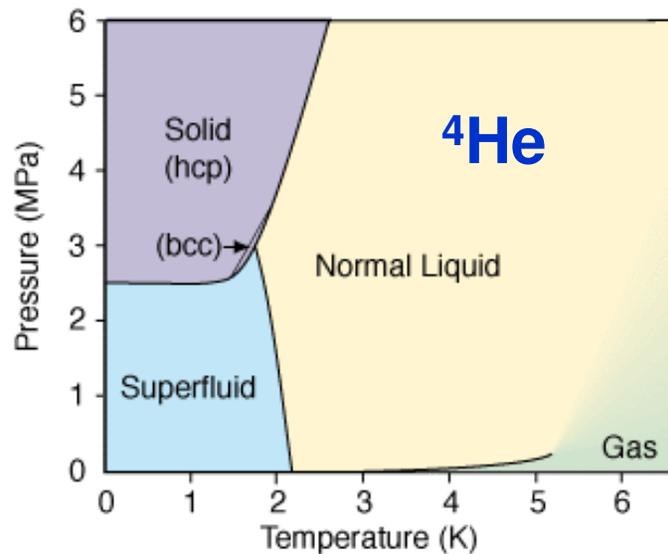
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Transformation from liquid to gas without a phase transition requires a process with pressures above the critical pressure.

There is only one phase of fluid above the critical pressure.



## Phase Transitions in Helium



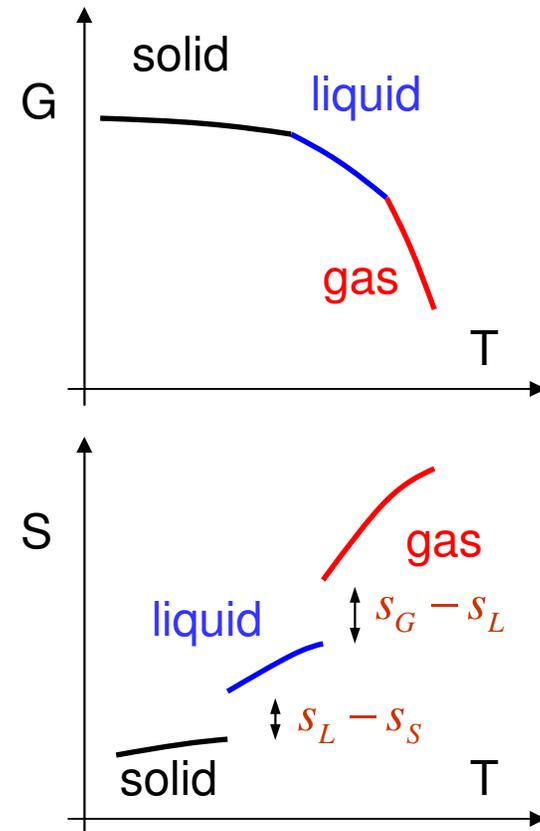
The vdW model of a fluid fails at low temperatures where interaction energy between atoms and molecules becomes comparable to temperature.

- Quantum phases such as superfluid phases of He are not described (a fluid with zero viscosity)

## Summary of phase transitions so far

- Entropy is discontinuous across the phase coexistence curve, e. g. molar entropies of liquid and gas phases at the same pressure and temperature are different. This also implies that a system undergoing a first order phase transition absorbs or emits heat at constant temperature (latent heat).
- Thermodynamically stable isotherms can be constructed from the unstable isotherms by using the fact the pressure and chemical potential remain constant across the phase transition

$$dG = -SdT + VdP + \mu dN \quad \Rightarrow \quad S = -\left(\frac{\partial G}{\partial T}\right)_{P,N}$$



## *First order phase transitions in multi-component systems*

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A two-component system where each of the components can be either solid or liquid.

$\mu_1^L(T, P, x_1^L)$  - chemical potential of the first component in the liquid state

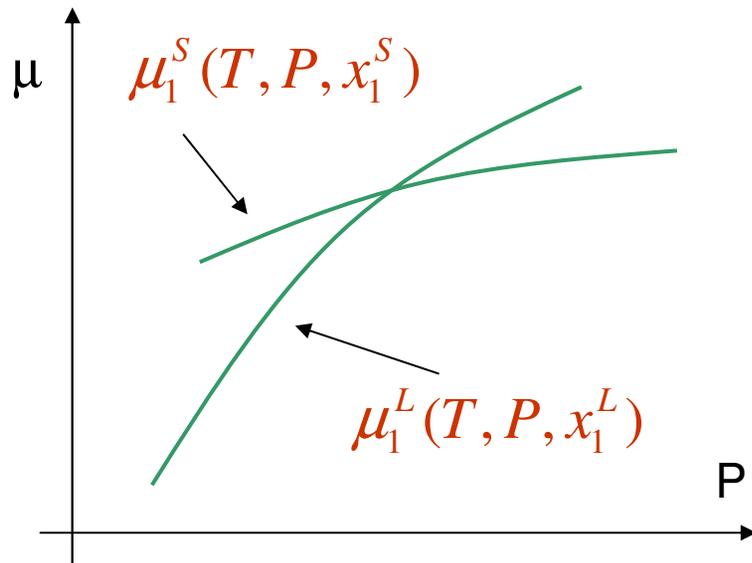
$\mu_1^S(T, P, x_1^S)$  - chemical potential of the first component in the solid state

$x_1^L$  - molar fraction of the first component in the liquid phase  $x_1^L + x_2^L = 1$

$x_1^S$  - molar fraction of the first component in the solid phase  $x_1^S + x_2^S = 1$

## Two-phase, two-component system

For component 1:



A solution exists in a 2D region of the P-T phase diagram.

The liquid and solid phases do coexist at the point when chemical potentials of the two phases are equal to each other:

$$\mu_1^S(T, P, x_1^S) = \mu_1^L(T, P, x_1^L) \quad (1)$$

$$\mu_2^S(T, P, x_2^S) = \mu_2^L(T, P, x_2^L)$$



$$\mu_2^S(T, P, 1 - x_1^S) = \mu_2^L(T, P, 1 - x_1^L) \quad (2)$$

Solving (1) and (2), we find

←  $x_1^L(P, T)$  and  $x_1^S(P, T)$

## *Three-phase, two-component system*

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Now assume that both components can exist in three phases: solid, liquid and gas

Gas, liquid and solid for both phases coexist if:

$$\mu_1^S(T, P, x_1^S) = \mu_1^L(T, P, x_1^L) = \mu_1^G(T, P, x_1^G)$$

$$\mu_2^S(T, P, 1 - x_1^S) = \mu_2^L(T, P, 1 - x_1^L) = \mu_2^G(T, P, 1 - x_1^G)$$

Note that these are four equations for three unknowns:  $x_1^S$ ,  $x_1^L$  and  $x_1^G$

## Three-phase, two-component system

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These are four equations for three unknowns:  $x_1^S$ ,  $x_1^L$  and  $x_1^G$

These means that three phases cannot coexist at arbitrary values of  $P$  and  $T$ . For a given value of  $T$ , the above 4 equations give  $P$ ,  $x_1^S$ ,  $x_1^L$  and  $x_1^G$

Similarly, if four phases of a two-component system are possible, they can only coexist at a uniquely defined point (or a few points) with given pressure and temperature.

Five phases cannot generally coexist in a two-component system.

# Gibbs Phase Rule

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For an arbitrary system with  $r$  components and  $M$  phases.

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## 1-component system:

1 phase: exists at any T and P	2 degrees of freedom
2 phases: coexist on the coexistence curves T(P)	1 degree of freedom
3 phases: coexist at a single point $T_t, P_t$	0 degrees of freedom

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## 2-component system:

2 phases: coexist in 2D regions of the T-P plane	2 degrees of freedom
3 phases: coexist in 1D regions T(P) of the T-P plane	1 degree of freedom
4 phases: coexist at a finite set of points $\{T_q, P_q\}$	0 degrees of freedom

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## r-component system:

$M$  phases:  $2+r-M$  degrees of freedom

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## Gibbs Phase Rule

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Gibbs phase rule: for  $r$ -component system,  $M$  phases can coexist in  $2+r-M$  dimensional regions of the thermodynamic coordinate space.

$$g = g(T, P, x_1, \dots, x_r) \quad x_1 + \dots + x_r = 1$$

:  $r+1$  thermodynamic coordinates

$$M \text{ phases coexist: } \mu_1(T, P, x_1, \dots, x_r) = \dots = \mu_M(T, P, x_1, \dots, x_r)$$

:  $M-1$  equations connecting thermodynamic coordinates

Therefore, there are  $r+1-(M-1) = 2+r-M$  independent thermodynamic coordinates in the regions of thermodynamic coordinate space where  $M$  phases coexist.