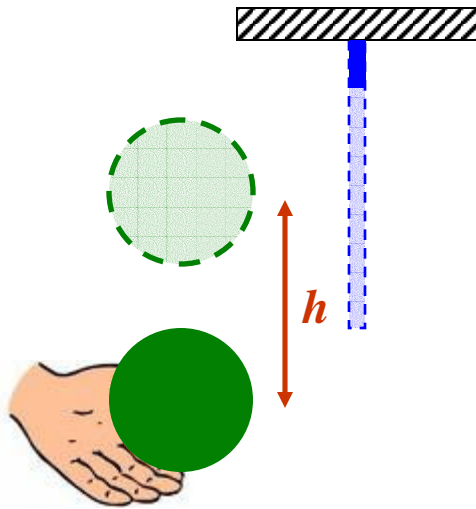


Energy-efficient use of Entropy

Problem: to minimize energy E spent to lift a weight using a stretched rubber band.

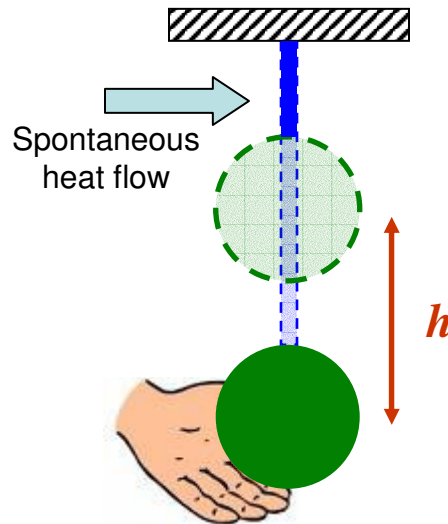
Lift by hand, no help from the rubber band

$$E_0 = W = Mgh$$



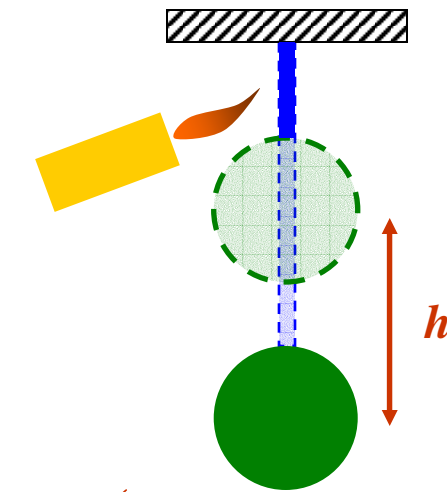
Lift by hand at constant T , rubber band helps

$$E_{T=const} = W' = ?$$



Lift by warming the rubber band

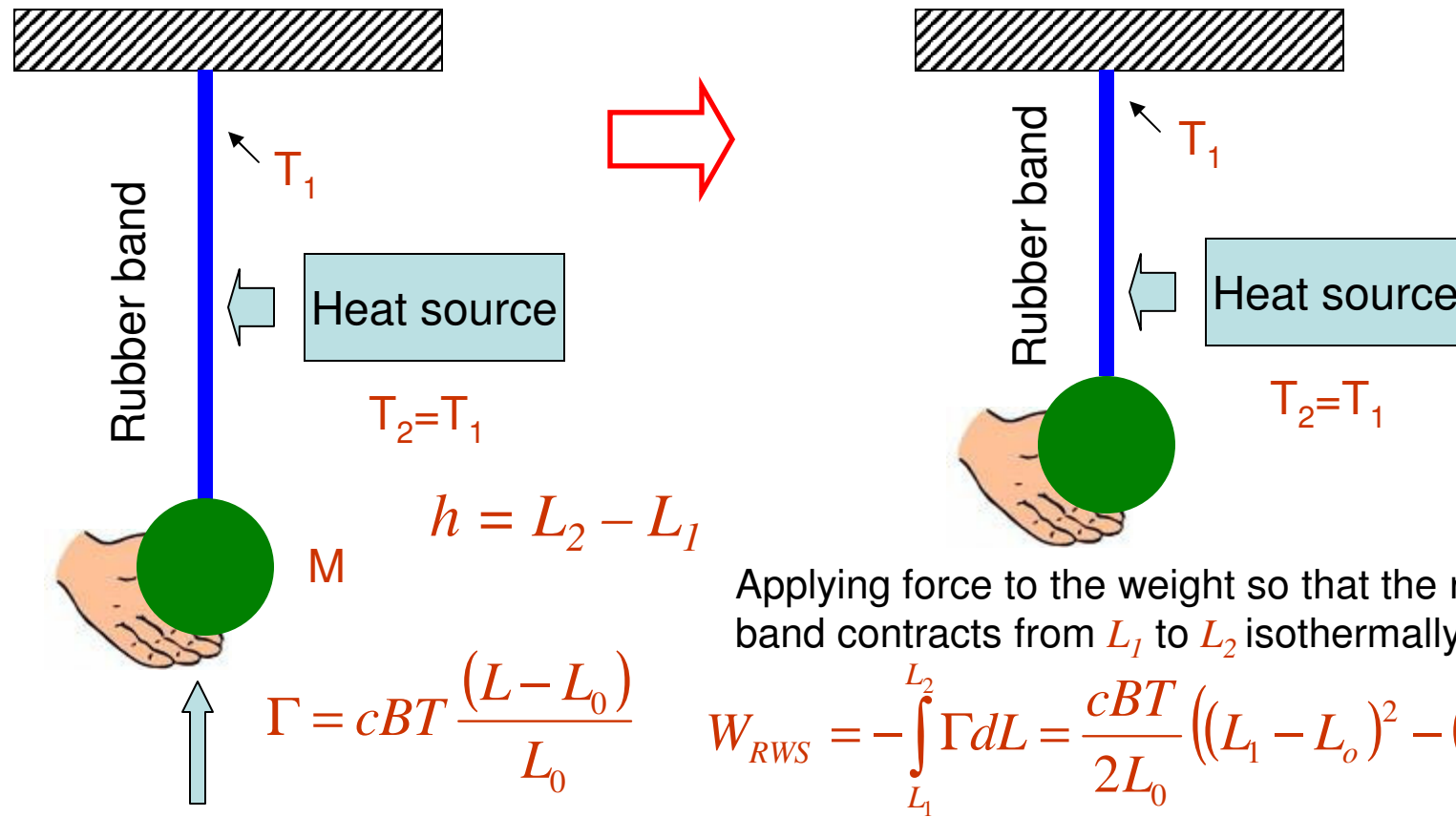
$$E_{T=const} = Q = ?$$



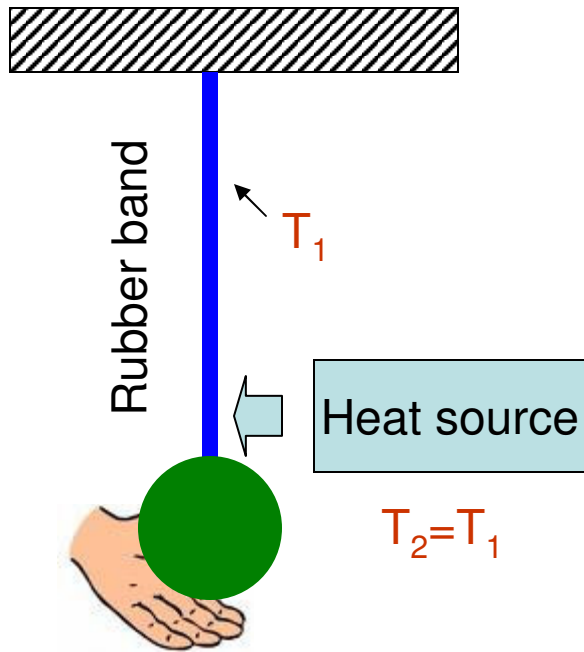
$$\min(E_{T=const}, E_{T=const}) = ?$$

Notes

Example of a heat engine: isothermal rubber band



Example of a heat engine: isothermal rubber band



Reversible heat engine

$$U = cL_0T \rightarrow \Delta U = 0$$

$$\Delta U = \Delta Q_{RHS} + \Delta W_{RWS} \rightarrow \Delta Q_{RHS} = -\Delta W_{RWS}$$

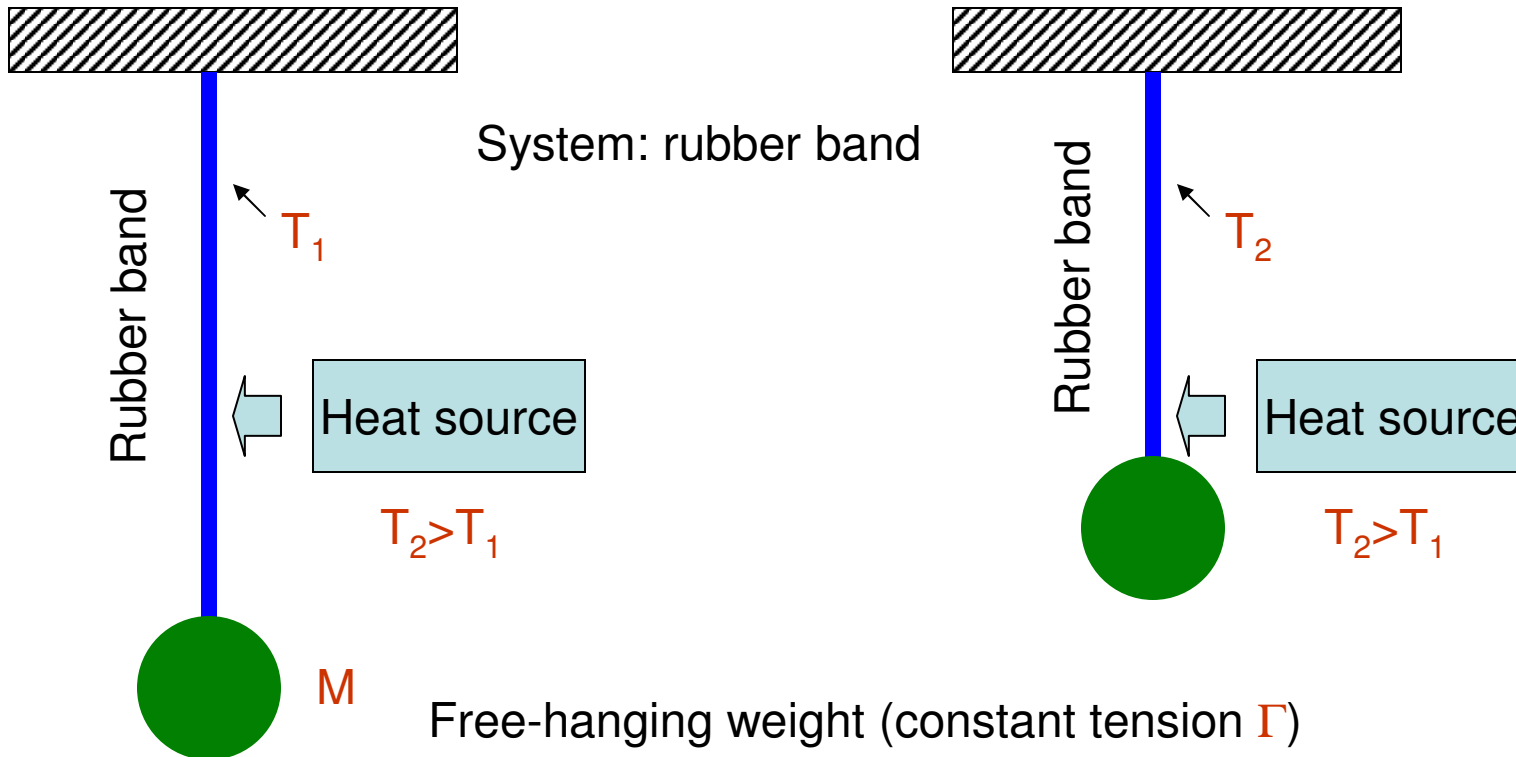
$$S_{band} = cL_0 \ln(U) - \frac{cB}{2} \frac{(L-L_0)^2}{L_0} + const$$

$$\Delta S_{band} = \frac{\Delta W_{RWS}}{T}$$

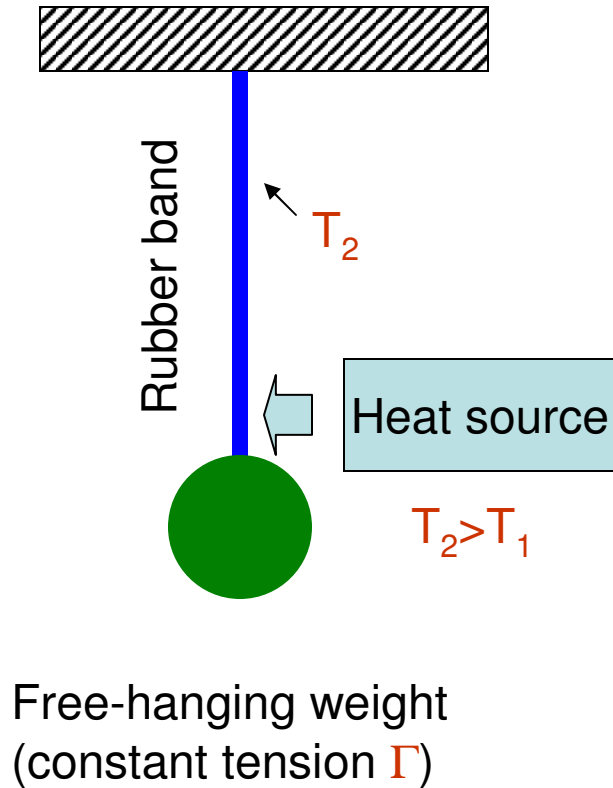
$$\Delta S_{RHS} = \frac{\Delta Q_{RHS}}{T}$$

$$\Delta S_{total} = \Delta S_{band} + \Delta S_{RHS} \rightarrow \Delta S_{total} = 0$$

Example of a heat engine: isobaric rubber band



Example of a heat engine: isobaric rubber band



$$U = cL_0T \quad \Rightarrow \quad \Delta U = cL_0(T_2 - T_1)$$

$$\Gamma = cBT \frac{(L - L_0)}{L_0} = Mg$$



$$h = L_2 - L_1 = \frac{MgL_0}{cB} \left(\frac{1}{T_2} - \frac{1}{T_1} \right)$$

Example of a heat engine: isobaric rubber band

$$\Delta U = cL_0(T_2 - T_1)$$

Work done by the rubber band

$$\Delta W_{RWS} = Mgh = \frac{M^2 g^2 L_0 (T_2 - T_1)}{cBT_1T_2} > 0$$

Heat transferred from the rubber band

$$\Delta Q_{RHS} = -\Delta U - \Delta W_{RWS} = -cL_0(T_2 - T_1) - Mgh = -cL_0(T_2 - T_1) \left(1 + \frac{M^2 g^2}{c^2 BT_1T_2} \right) < 0$$

Heat transferred to the rubber band

$$E_{\Gamma=const} = Q = |\Delta Q_{RHS}| = cL_0(T_2 - T_1) + Mgh > Mgh$$

Example of a heat engine: isobaric rubber band

$$S = cL_0 \ln(U) - \frac{cB}{2} \frac{(L - L_0)^2}{L_0} + \text{const}$$

Change of entropy of the band:

$$\begin{aligned}\Delta S &= cL_0 \ln\left(\frac{T_2}{T_1}\right) - \frac{cB}{2L_0} \left((L_2 - L_0)^2 - (L_1 - L_0)^2 \right) = \\ &= cL_0 \ln\left(\frac{T_2}{T_1}\right) + \frac{cB}{2L_0} \left(\frac{MgL_0}{cB} \right)^2 \frac{(T_2 - T_1)(T_2 + T_1)}{T_1^2 T_2^2}\end{aligned}$$

$$\Delta S = cL_0 \ln\left(\frac{T_2}{T_1}\right) + \frac{cB}{2L_0 T_2^2} \left(\frac{MgL_0}{cB} \right)^2 \left(\frac{T_2^2}{T_1^2} - 1 \right)$$

Example of a heat engine: isobaric rubber band

Change of entropy of the heat source:

$$\Delta S_{RHS} = \frac{\Delta Q_{RHS}}{T_2} = -cL_0 \left(\frac{T_2 - T_1}{T_2} \right) \left(1 + \frac{M^2 g^2}{c^2 B T_1 T_2} \right)$$

Total entropy change of the system: $\Delta S_{total} = \Delta S_{RHS} + \Delta S$

$$\Delta S_{total} = -cL_0 \frac{(T_2 - T_1)}{T_2} \left(1 + \frac{M^2 g^2}{c^2 B T_1 T_2} \right) + cL_0 \ln \left(\frac{T_2}{T_1} \right) + \frac{cB}{2L_0 T_2^2} \left(\frac{MgL_0}{cB} \right)^2 \left(\frac{T_2^2}{T_1^2} - 1 \right)$$

$$\Delta S_{total} = cL_0 \ln \left(\frac{T_2}{T_1} \right) - cL_0 \left(1 - \frac{T_1}{T_2} \right) + \frac{M^2 g^2 L_0}{2cB T_2^2} \left(\frac{T_2}{T_1} - 1 \right)^2$$

Example of a heat engine: isobaric rubber band

$$\Delta S_{total} = cL_0 \ln\left(\frac{T_2}{T_1}\right) - cL_0\left(1 - \frac{T_1}{T_2}\right) + \frac{M^2 g^2 L_0}{2cBT_2^2} \left(\frac{T_2}{T_1} - 1\right)^2$$

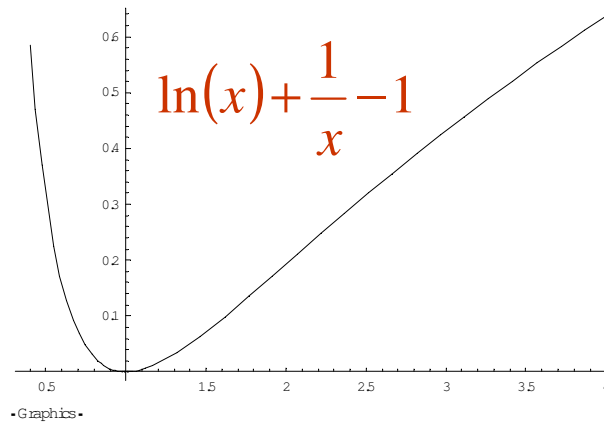
$$\Delta S_{total} \geq 0 \quad \text{because} \quad \frac{M^2 g^2 L_0}{2cBT_2^2} \left(\frac{T_2}{T_1} - 1\right)^2 \geq 0 \quad \text{and} \quad cL_0 \ln\left(\frac{T_2}{T_1}\right) - cL_0\left(1 - \frac{T_1}{T_2}\right) \geq 0$$

The latter is equivalent to:

$$\ln(x) + \frac{1}{x} \geq 1$$

Which is a true statement

Plot[Log[x] + 1/x - 1, {x, 0.4, 4}, PlotRange -> All]



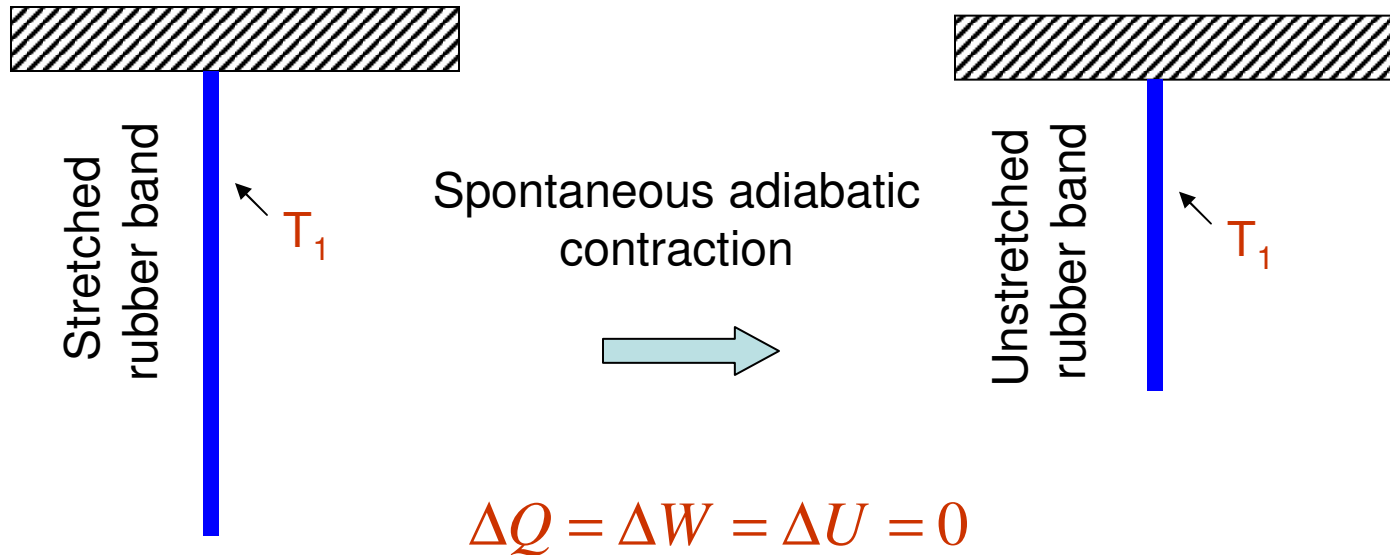
$$\Delta S_{total} > 0$$

If $T_2 \neq T_1$

Isobaric rubber band engine is irreversible!

Harvesting Low Entropy to do work: rubber band

Case 1 – low entropy wasted



Harvesting Low Entropy to do work: rubber band

$$\Delta Q = \Delta W = \Delta U = 0$$

Initial entropy

final entropy

$$S_i = cL_0 \ln(U) - \frac{cB}{2} \frac{(L - L_0)^2}{L_0} + \text{const}$$

$$S_f = cL_0 \ln(U) + \text{const}$$

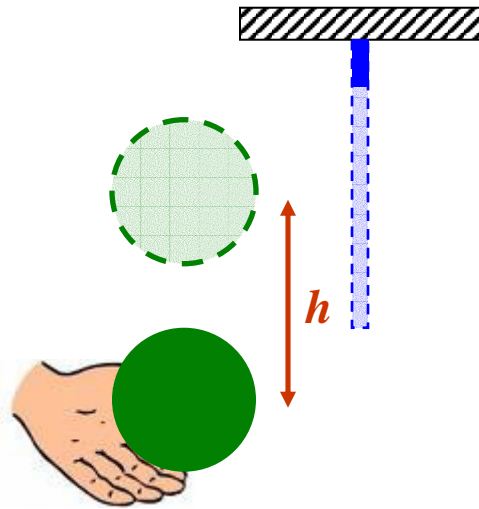
$$\Delta S_{total} = \Delta S = S_f - S_i = \frac{cB}{2} \frac{(L - L_0)^2}{L_0} > 0$$

No work done, total entropy increased

Summary

Lift by hand, no help from the rubber band

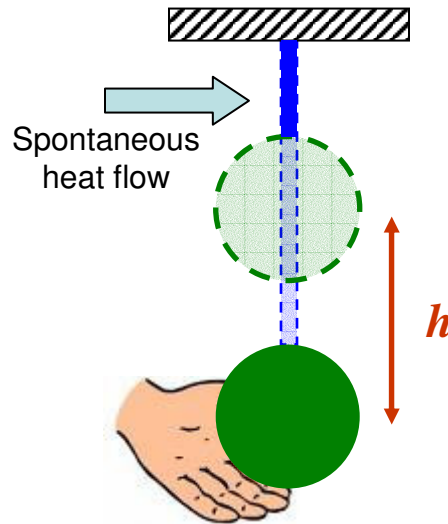
$$E_0 = W = Mgh$$



$$\Delta S_{total} > 0$$

Lift by hand at constant T , rubber band helps

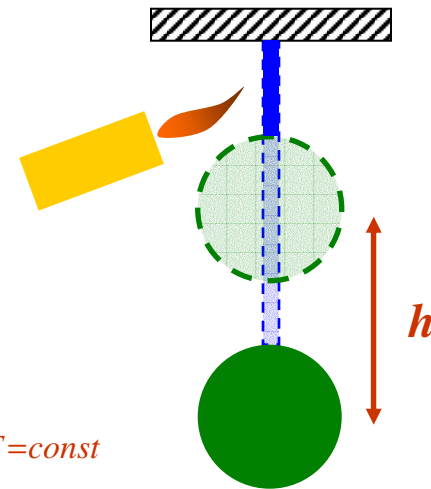
$$E_{T=const} = W < Mgh$$



$$\Delta S_{total} = 0$$

Lift by warming the rubber band

$$E_{T=const} = Q = Mgh + \Delta U$$



$$\Delta S_{total} > 0$$

$$E_{min} = E_{T=const}$$

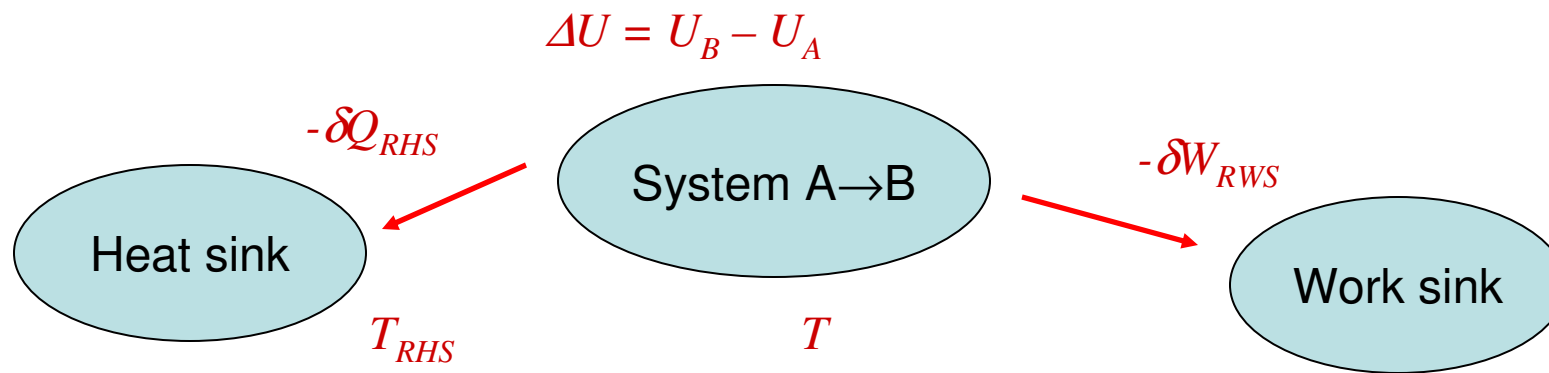
Least energy is spent in the reversible process to lift the weight

The Maximum Work Theorem

We model the environment of the system as two objects:

- reversible work source (sink) RWS
- reversible heat source (sink) RHS

Heat engine model



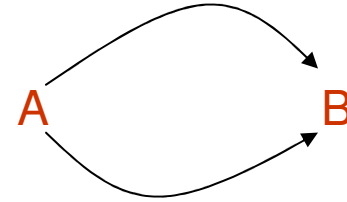
Rigid walls, so no work can be done on the heat sink !

Adiabatic walls, so no heat can be transferred to the work sink! ($dS_{RWS}=0$)

The Maximum Work Theorem

Maximum work theorem: for a given ΔU , $-\delta W$ is maximum for a reversible process

System goes from state A to state B.
Which process delivers maximum work
to the reversible work source?



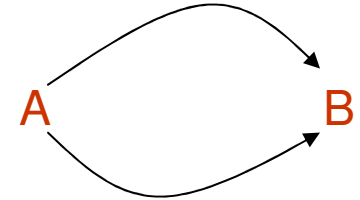
The change of internal energy ΔU is the same for all A→B processes

ΔU is distributed between work done on RWS, ΔW , and heat flow to the RHS, ΔQ

Therefore, maximum work delivered to RWS implies minimum heat flow to RHS.

The Maximum Work Theorem

Maximum work theorem: for a given ΔU , $-\delta W$ is maximum for a reversible process



The change of entropy of the working body ΔS is the same for all A→B processes

The change of entropy of the RHS is proportional to ΔQ , therefore maximum work corresponds to the **minimum entropy change** of the RHS and (since ΔS is independent on the process A→B) of the entire composite system (engine).

The absolute minimum entropy increase is zero and **achieved in a reversible process**.

Therefore, maximum work done by a heat engine is achieved in a reversible process.

The Maximum Work Theorem

Maximum work theorem: for a given ΔU , $-\delta W$ is maximum for a reversible process

Proof:

$$dS_{total} = dS + dS_{RHS} = dS + \frac{\delta Q_{RHS}}{T_{RHS}} \geq 0 \quad \Rightarrow \quad \delta Q_{RHS} \geq -T_{RHS} dS$$

$$dU = -\delta W_{RWS} - \delta Q_{RHS}$$



$$dU \leq -\delta W_{RWS} + T_{RHS} dS$$



$$\boxed{\delta W_{RWS} \leq T_{RHS} dS - dU}$$

Functions of states **A** and **B**
only, not of the process **A**→**B**

The Maximum Work Theorem

$$\delta W_{RWS} \leq T_{RHS} dS - dU$$

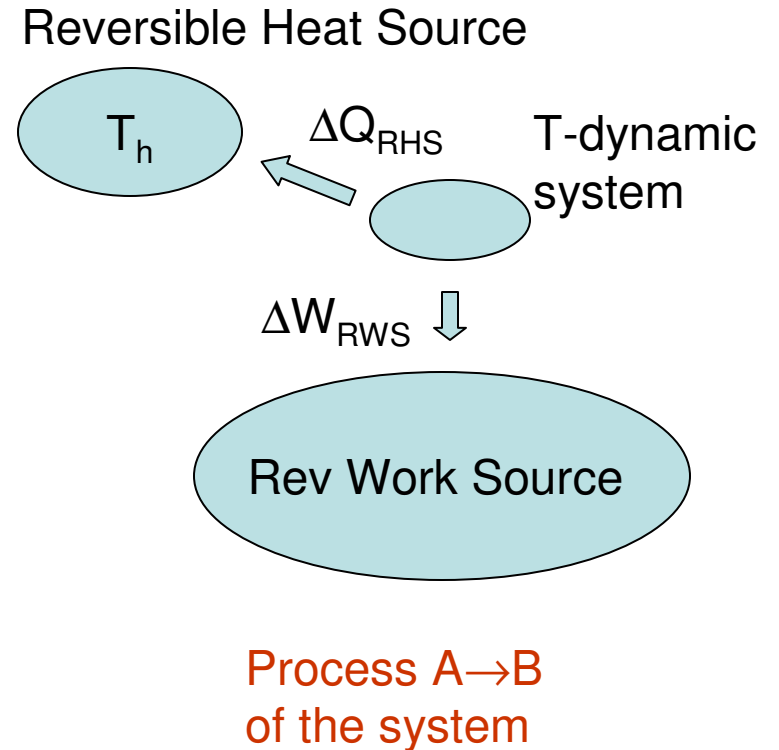
The work is maximum when $\leq \rightarrow =$

$$\delta W_{RWS}^{\max} = T_{RHS} dS - dU$$

The origin of the $\leq$ sign is traced back to $dS_{total} \geq 0$.
Therefore, for maximum work we have $dS_{total} = 0$,
hence *reversible process maximizes work*.

Summary of heat engines so far

- A thermodynamic system in a process connecting state **A** to state **B**
- In this process, the system can do work and emit/absorb heat
- What processes maximize work done by the system?
- We have proven that **reversible processes** (processes that do not change the total entropy of the system and its environment) **maximize work** done by the system in the process **A→B**



Cyclic Engine

-To **continuously generate work**, a heat engine should go through a **cyclic process**.

- Such an engine needs to have a heat source and a heat sink at different temperatures.

- The net result of each cycle is to extract heat from the heat source and distribute it between the heat transferred to the sink and useful work.

